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Uniqueness in nonlinearly coupled PDE systems
(joint work with Lucia Panizzi)

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Model problem: surface hardening of steel

In a domain $Q_T = \Omega \times (0, T)$, $\Omega \subset \mathbb{R}^3$ bounded and smooth, consider the system

$$\begin{aligned}\theta_t - \Delta\theta &= r(\theta, c) \\ c_t - \operatorname{div}(D(\theta, c)\nabla c) &= 0\end{aligned}$$

with boundary conditions on $\Sigma_T = \partial\Omega \times (0, T)$

$$\begin{aligned}\frac{\partial\theta}{\partial\nu} + h(x, \theta, \theta_\Gamma(x, t)) &= 0 \\ -D(\theta, c)\frac{\partial c}{\partial\nu} &= b(x, t)\end{aligned}$$

and initial conditions $\theta(x, 0) = \theta^0(x)$, $c(x, 0) = c^0(x)$.

$\theta^0 \geq \theta_* > 0$, $\theta_\Gamma \geq \theta_*$, $b \in L^2(\Sigma_T)$, $c^0 \in L^2(\Omega)$... given data

$r \geq 0$, $D \in [d_0, d_1]$, $h(x, \cdot, \cdot)$... (smooth) constitutive functions

Existence and regularity I.

Assumption on h :

$$\exists a > 0 : \theta, \theta_\Gamma \geq \theta_* \implies h(x, \theta, \theta_\Gamma)(\theta - \theta_\Gamma) \geq a(\theta - \theta_\Gamma)^2.$$

Existence can be proved by Schauder argument.

Maximal expected regularity for c : $\nabla c \in L^2(Q_T)$

$$c_t \in L^2(0, T; (H^1(\Omega))')$$

$$c \in L^\infty(0, T; L^2(\Omega))$$

Regularity for θ : $\theta_t, \Delta \theta \in L^2(Q_T)$

$$\theta \geq \theta_* \quad (\text{maximum principle})$$

$$\theta \in L^\infty(Q_T) \quad (\text{Alikakos-Moser})$$

This is not enough for uniqueness !

Partial Kirchhoff transform

We introduce new functions

$$u = F(\theta, c) := \int_0^c D(\theta, c') dc' \iff c = G(\theta, u)$$

$$R(\theta, u) := r(\theta, G(\theta, u)) , \quad H(\theta, u) := \partial_\theta F(\theta, G(\theta, u)) , \quad U(x, t) := \int_0^t u(x, \tau) d\tau$$

We assume that F, G, R, H are Lipschitz, and H is bounded.

In terms of θ, U , the system of equations has the form

$$\int_{\Omega} (\theta_t \varphi + \nabla \theta \cdot \nabla \varphi - R(\theta, U_t) \varphi) dx + \int_{\partial \Omega} h(x, \theta, \theta_\Gamma(x, t)) \varphi dS = 0$$

$$\int_{\Omega} (G(\theta, U_t)_t \psi + \nabla U_t \cdot \nabla \psi - H(\theta, U_t) \nabla \theta \cdot \nabla \psi) dx + \int_{\partial \Omega} b(x, t) \psi dS = 0$$

for every test functions $\varphi, \psi \in H^1(\Omega)$.

Uniqueness I.

Consider two solutions (θ_1, U_1) , (θ_2, U_2) , and set $\bar{\theta} = \theta_1 - \theta_2$, $\bar{U} = U_1 - U_2$.

Test the difference of the equations for θ_1 , θ_2 by $\varphi = \bar{\theta}$;

Integrate the equations for U_1 , U_2 in time, and test the difference by $\psi = \bar{U}_t$.

On the left hand side, we have

$$\frac{d}{dt} \left(|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 \right) + |\nabla \bar{\theta}(t)|_{L^2(\Omega)}^2 + |\bar{U}_t(t)|_{L^2(\Omega)}^2$$

We omit the good terms on the right hand side that can be estimated by Gronwall's argument. The only bad term is the integral

$$\int_{\Omega} \left(\int_0^t (H(\theta_1, (U_1)_t) \nabla \theta_1 - H(\theta_2, (U_2)_t) \nabla \theta_2) d\tau \right) \cdot \nabla \bar{U}_t(x, t) dx$$

We have to integrate by parts in time !

Uniqueness II.

Two bad terms remain:

$$(1) \frac{d}{dt} \int_{\Omega} \left(\int_0^t (H(\theta_1, (U_1)_t) \nabla \theta_1 - H(\theta_2, (U_2)_t) \nabla \theta_2) d\tau \right) \cdot \nabla \bar{U}(x, t) dx$$

$$(2) \int_{\Omega} (H(\theta_1, (U_1)_t) \nabla \theta_1 - H(\theta_2, (U_2)_t) \nabla \theta_2) \cdot \nabla \bar{U}(x, t) dx$$

We integrate in time. Omitting again the terms that can be handled via Gronwall, we obtain for the left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

the upper bounds

$$(1)^* \int_{\Omega} |\nabla \bar{U}(x, t)| \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right) dx$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau$$

Regularity II.

A necessary and sufficient condition to continue the above uniqueness estimates reads

$$\nabla \theta \in L^2(0, T; L^\infty(\Omega))$$

It is proved by differentiating θ in tangential directions and using the boundary condition to check the regularity of the normal component.

- Boundary of Ω must be of class $C^{2,1}$
- Different regularity is obtained for the normal and for the tangential derivatives
- Anisotropic Sobolev embedding theorem

For a vector $\mathbf{p} = (p_1, \dots, p_N)$ we define the space $L^\mathbf{p}(\mathbb{R}^N)$ as the subspace of $L^1(\mathbb{R}^N)$ of functions u with finite norm

$$\|u\|_{\mathbf{p}} = \left(\int_{\mathbb{R}} \left(\dots \int_{\mathbb{R}} \left(\int_{\mathbb{R}} |u(x)|^{p_1} dx_1 \right)^{p_2/p_1} dx_2 \dots \right)^{p_N/p_{N-1}} dx_N \right)^{1/p_N}$$

Regularity III.

For a matrix $\mathbf{P} = (P_{ij})_{i,j=1}^N$, $P_{ij} = 1/p_{ij}$, we define the anisotropic Sobolev space

$$W^{1,\mathbf{P}}(\mathbb{R}^N) = \left\{ u \in L^1(\mathbb{R}^N) : \frac{\partial u}{\partial x_i} \in L^{\mathbf{p}_i}(\mathbb{R}^N), i = 1, \dots, N \right\}, \quad \mathbf{p}_i = (p_{i1}, \dots, p_{iN})$$

We denote by $\varrho(\mathbf{P})$ the spectral radius of \mathbf{P} , by \mathbf{I} the identity $N \times N$ matrix, and by $\mathbf{1}$ the vector $\mathbf{1} = (1, 1, \dots, 1)$.

Theorem Let $\varrho(\mathbf{P}) < 1$, and let

$$(\mathbf{I} - \mathbf{P})^{-1} \mathbf{1} = \mathbf{b} = (b_1, \dots, b_N).$$

Then $W^{1,\mathbf{P}}(\mathbb{R}^N)$ is compactly embedded in $L^\infty(\mathbb{R}^N)$, and there exists a constant $C > 0$ such that each $u \in W^{1,\mathbf{P}}(\mathbb{R}^N)$ has for all $x, z \in \mathbb{R}^N$ the Hölder property

$$|u(z) - u(x)| \leq C \|u\|_{W^{1,\mathbf{P}}(\mathbb{R}^N)} \sum_{i=1}^N |z_i - x_i|^{1/b_i}.$$

Uniqueness III.

In the uniqueness estimates with left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

and right hand side

$$(1)^* \int_{\Omega} |\nabla \bar{U}(x, t)| \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right) dx$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau$$

we first use the Cauchy-Schwarz inequality

Uniqueness III.

In the uniqueness estimates with left hand side

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and right hand side

$$(1)^* |\nabla \bar{U}(t)|_{L^2(\Omega)} \left(\int_{\Omega} \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx \right)^{1/2},$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau,$$

we first use the Cauchy-Schwarz inequality

Uniqueness III.

In the uniqueness estimates with left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

and right hand side

$$(1)^* \int_{\Omega} \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx ,$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau ,$$

then the Young inequality

Uniqueness III.

In the uniqueness estimates with left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

and right hand side

$$(1)^* \int_{\Omega} \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx ,$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau ,$$

we finally eliminate again the terms subject to **Gronwall's argument**.

Uniqueness III.

In the uniqueness estimates with left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

and right hand side

$$(1)^* \int_{\Omega} \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx ,$$

$$(2)^* \int_0^t \int_{\Omega} |\nabla \bar{U}| (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) dx d\tau ,$$

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and right hand side

$$(1)^* \int_{\Omega} \left(\int_0^t (|\bar{\theta}| + |\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx ,$$

we finally eliminate again the terms subject to **Gronwall's argument**.

Uniqueness III.

In the uniqueness estimates with left hand side

$$|\bar{\theta}(t)|_{L^2(\Omega)}^2 + |\nabla \bar{U}(t)|_{L^2(\Omega)}^2 + \int_0^t \left(|\nabla \bar{\theta}(\tau)|_{L^2(\Omega)}^2 + |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 \right) d\tau$$

and right hand side

$$(1)^* \int_{\Omega} \left(\int_0^t (|\bar{U}_t|) |\nabla \theta_1|(x, \tau) d\tau \right)^2 dx ,$$

the remaining term is estimated using Minkowski's inequality

$$\int_{\Omega} \left(\int_0^t |\bar{U}_t| |\nabla \theta_1| d\tau \right)^2 dx \leq \left(\int_0^t \left(\int_{\Omega} |\bar{U}_t|^2 |\nabla \theta_1|^2 dx \right)^{1/2} d\tau \right)^2$$

Uniqueness IV.

The critical inequality has the form

$$\int_0^t |\bar{U}_t(\tau)|_{L^2(\Omega)}^2 d\tau \leq \left(\int_0^t \gamma(\tau) |\bar{U}_t(\tau)|_{L^2(\Omega)} d\tau \right)^2$$

with a function $\gamma \in L^2(0, T)$: $\boxed{\gamma(t) = C \operatorname{ess\ sup}_{x \in \Omega} |\nabla \theta_1(x, t)|}$.

Lemma (L^p Gronwall inequality)

Let $s \in L^p(0, T)$, $\gamma \in L^{p'}(0, T)$, $\alpha \in L^\infty(0, T)$ be nonnegative functions such that

$$\left(\int_0^t s^p(\tau) d\tau \right)^{1/p} \leq \alpha(t) + \int_0^t \gamma(\tau) s(\tau) d\tau$$

for every $t \in [0, T]$, with a fixed $p > 1$. Then there exists $M_\gamma > 0$ such that

$$\left(\int_0^T s^p(\tau) d\tau \right)^{1/p} \leq M_\gamma \operatorname{ess\ sup}_{[0, T]} \alpha.$$

Conclusions

- In coupled parabolic systems, lack of regularity in one equation can be compensated by a higher regularity in the other equation;
- Anisotropic Sobolev embeddings might be useful;
- Gronwall's inequality is not the only argument for uniqueness;
- These observations might be of interest also in other situations (Penrose-Fife with state dependent heat conductivity, say; In this case, the phase variable is the more regular one).