

Convergence to equilibria of solutions of nonlocal phase-separation models

$$\vartheta_t + \chi_t - \Delta \vartheta = f$$

$$\partial_t \chi + \chi - J * \chi + W'(\chi) = \vartheta, \quad t > 0, \quad x \in \Omega \subset R^3,$$

$$\varepsilon \partial_{tt} \chi + \partial_t \chi + \chi - J * \chi + W'(\chi) = \vartheta,$$

where

$$J * \chi(x, t) = \int_{\Omega} J(x - y) \chi(y, t) \, dy \stackrel{def}{=} \mathcal{J}[\chi],$$

$$u(x, 0) = u_0(x) \in L^{\infty}(\Omega)$$

$$u_t - \nabla \cdot (\mu \nabla v) = 0$$

$$v = f'(u) + \int_{\Omega} J(|x - y|)(1 - 2u(t, y)) \, dy,$$

$$\mu \nu \cdot \nabla v = 0 \text{ in } (0, T) \times \partial \Omega,$$

$$f(u) = u \ln u + (1 - u) \ln(1 - u), \quad \mu = \frac{a}{f''(u)},$$

$$u_0 \in L^{\infty}(\Omega), \quad 0 \leq u_0(x) \leq 1, \quad 0 < \int_{\Omega} u_0 = u_{\alpha} < 1.$$

$$\dot{u}(t) + M(u(t)) = g(t), \quad u(0) = u_0 \in H$$

$$\ddot{u}(t) + \dot{u}(t) + M(u(t)) = g(t), \quad u(0) = u_0 \in V, \dot{u}(0) = u_1 \in H$$

$$V \subset H \subset V', \quad M = D E$$

$$|E(v) - E(\phi)|^{1-\theta} \leq C \|M(v)\|_{V'}, \quad \phi \in \omega(u)$$

$\{u(t); t \geq 0\}$ is relatively compact in V

Then

$$\lim_{t \rightarrow \infty} \|u(t) - \phi\|_V = 0$$

The energy functional takes a form

$$E[u] = \int_{\Omega} \left(\frac{1}{2} |\nabla u|^2 + F(u) \right) dx,$$

i.e., it is a quadratic form perturbed by a non-linear compact functional. The underlying function spaces are the Sobolev spaces $W^{1,p}(\Omega)$ on which E is analytic provided p is large enough and F is an analytic function.

The energy functional related to our problems reads

$$E[\chi] = E[\chi, 0] = \int_{\Omega} \left(\frac{1}{2} \chi(\chi - \mathcal{J}[\chi]) + W(\chi) \right) dx.$$

the natural domain of E is $L^2(\Omega)$.

$$E \in C^1(L^2(\Omega); R), \quad E \notin C^2(L^2(\Omega); R)$$

no matter how smooth W is.

$$u_t - \nabla \cdot (\mu \nabla v) = 0$$

$$v = f'(u) + \int_{\Omega} J(|x - y|)(1 - 2u(t, y))dy = \log \frac{u}{1 - u} + w,$$

$$\mu \nu \cdot \nabla v = 0 \text{ in } (0, T) \times \partial\Omega,$$

$$f(u) = u \ln u + (1 - u) \ln(1 - u), \quad \mu = \frac{a}{f''(u)} = u(1 - u),$$

$$u_0 \in L^{\infty}(\Omega), \quad 0 \leq u_0(x) \leq 1, \quad 0 < \int_{\Omega} u_0 = u_{\alpha} < 1.$$

Stationary solutions

$$u^* = \frac{1}{1 + \exp(w^* - v^*)}, \quad v^* = \text{const},$$

$$w^* = \int_{\Omega} K(|x - y|)(1 - 2u^*(t, y))dy$$

The energy functional

$$F(z) = \int_{\Omega} f(z) + u\mathcal{J}(z) + z \cdot K * 1$$

$$\frac{d}{dt}F(u(t)) = - \int_{\Omega} \mu |\nabla v|^2 dx$$

$$F'(u(t)) = f'(u(t)) - 2\mathcal{J}(u(t)) + l = v(t)$$

$$\int_t^{\infty} \int_{\Omega} \mu |\nabla v|^2 dx ds = F(u(t)) - F_{\infty} = F(u(t)) - F(u^*).$$

By virtue of the generalized Lojasiewicz theorem (Gajewski, Griepentrog)

$$|F(u(t)) - F(u^*)|^{1-\theta} \leq$$

$$\lambda \inf\{\|v(t) - z\|_{L^2(\Omega)}; z = \text{const}\} = \lambda \|v(t) - \int_{\Omega} v(t)\|_{L^2(\Omega)}$$

for

$$\|u(t) - u^*\|_{L^2(\Omega)} \leq \delta.$$

$$\begin{aligned} \frac{4}{a} \int_t^\infty \int_{\Omega} (\mu |\nabla v|)^2 dx ds &\leq \int_t^\infty \int_{\Omega} \mu |\nabla v|^2 dx ds \\ &\leq \lambda \|v(t) - \int_{\Omega} v(t)\|_{L^2(\Omega)}^{\frac{1}{1-\theta}} \leq \|\mu |\nabla v|(t)\|_{L^2(\Omega)}^{\frac{1}{1-\theta}}, \end{aligned}$$

$$\Rightarrow u_t \in L^1(T, \infty; H^1(\Omega))^*$$

provided that $0 < k \leq \mu \leq 1 - k$ a.a. in $((T, \infty) \times \Omega)$

$$\|\log u_0\|_{L^\infty(\Omega)} < \infty$$

Separation property for general initial condition

$$u_0 \in L^\infty(\Omega), \quad 0 \leq u_0(x) \leq 1, \quad 0 < \int_{\Omega} u_0 = u_\alpha < 1.$$

$$\|\ln u(t)\|_{L^r(\Omega)} \leq B_1 \|\ln u(0)\|_{L^r(\Omega)}, \quad t \geq 0,$$

$$\|\ln u(t)\|_{L^r(\Omega)} \leq B_2 r^2, \quad t \geq 1,$$

$$\|\ln u(t)\|_{L^r(\Omega)} \leq B_3, \quad t \geq t_r.$$

Keller-Segel chemotaxis model

$$\partial_t u(t, x) = \operatorname{div}(\nabla u(t, x) - u(t, x)\nabla v(t, x)),$$

$$\partial_t v(t, x) = \alpha \Delta v(t, x) - \beta v(t, x) + \gamma(u(t, x) - 1) \text{ for } t > 0, x \in \Omega,$$

$$\nabla u(t, x) \cdot \nu(x) = \nabla v(t, x) \cdot \nu(x) = 0, \quad t > 0, x \in \partial\Omega,$$

$u = u(t, x)$ -the species density,

$v = v(t, x)$ - a rescaled chemical substance density.

Lyapunov function

$$F[u, v] \equiv \int_{\Omega} (u \log(u) + \frac{\alpha}{2\gamma} |\nabla v|^2 + \frac{\beta}{2\gamma} |v|^2 - uv),$$

defined on the space

$$X = \{(u, v) \mid v \in W^{1,2}(\Omega), u \in L^2(\Omega), \int_{\Omega} u = 0\}.$$

F is not twice continuously differentiable on X . The gradient ∂F is of type “monotone operator + linear compact perturbation” while the usual form is “linear isomorphism + (non-linear) compact perturbation”. The operator associated to the gradient ∂F in the sense of distribution does not map the space X into itself, it does not conserve the (zero) mean of the first component.

$$|F(u, v) - F(U, V)|^{1-\theta} \leq m \|\partial F(u, v)\|_{X^*}$$

$$\text{whenever } \|(u, v) - (U, V)\|_X < \epsilon.$$

U, V are solutions of the stationary problem.

References

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