

Asymptotics in Nonlinear Hyperbolic Thermoelasticity¹

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- Plan:
1. Introduction
 2. One-dimensional non-linear Cauchy problem
 3. Propagation of singularities in three dimensions
 4. Low frequency expansion in exterior domains

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1 Introduction

Thermoelastic systems for the displacement vector u and the the temperature difference θ , where the classical Fourier law

$$q + k\nabla\theta = 0$$

is replaced by Cattaneo's law

$$\tau q_t + q + k\nabla\theta = 0$$

τ : relaxation parameter.

$$\rho u_{tt} - \operatorname{div} S = 0 \quad (1.1)$$

$$\varepsilon_t - \operatorname{tr}\{S \nabla u_t\} + \operatorname{div} q = 0 \quad (1.2)$$

$$(\theta + T_0)\{a(\nabla u, \theta, q)\theta_t - \operatorname{tr}[S_\theta \cdot \nabla u_t]\} + \operatorname{div} q = b(\nabla u, \theta, q)q_t \quad (1.3)$$

Cattaneo's law:

$$\tau(\nabla u, \theta)q_t + q + k(\nabla u, \theta)\nabla \theta = 0 \quad (1.4)$$

Topics:

- Decay rates of solutions to the one-dimensional Cauchy problem
- Propagation of singularities depending on the curvature of the initial surface of discontinuities in three space dimensions
- Low frequency asymptotics in three-dimensional exterior domains as first step towards the time asymptotics in exterior domains

2 One-dimensional non-linear Cauchy problem

(Joint work with Y.-G. Wang)

Nonlinear Cauchy problem in one space dimension

$$u_{tt} - a_1(P)u_{xx} + a_2(P)\theta_x + \sigma_1(P)q_x = 0 \quad (2.1)$$

$$\theta_t + a_3(P)u_{tx} + a_4(P)q_x + \sigma_2(P)\theta_x + a_5(P)q = 0 \quad (2.2)$$

$$q_t + a_6(u_x, \theta)\theta_x + a_7(u_x, \theta)q = 0 \quad (2.3)$$

$$u(0, x) = u_0(x), \quad u_t(0, x) = u_1(x), \quad \theta(0, x) = \theta_0(x), \quad q(0, x) = q_0(x), \quad x \in \mathbb{R} \quad (2.4)$$

where

$$P = (u_x, \theta, q)$$

Transforming (2.1) - (2.4) into a first-order system for

$$V := (u_x, u_t, \theta, q)'$$

$$V_t + \underbrace{\begin{pmatrix} 0 & -1 & 0 & 0 \\ -a_1 & 0 & a_2 & \sigma_1 \\ 0 & a_3 & \sigma_2 & a_4 \\ 0 & 0 & a_6 & 0 \end{pmatrix}}_{=:A(V)} \partial_x V + \underbrace{\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & a_5 \\ 0 & 0 & 0 & a_7 \end{pmatrix}}_{=:B(V)} V = 0 \quad (2.5)$$

$$V(t=0) = V_0 := (u_{0,x}, u_1, \theta_0, q_0)' \quad (2.6)$$

This is strictly hyperbolic if $|P|$ is small enough, hence:

Theorem 2.1. *Let $s \geq 2$. Then there is $\delta > 0$ such that for data $V_0 \in W^{s,2}(\mathbb{R})$ with $\|V_0\|_{s,2} < \delta$ there is a unique local solution V to (2.5), (2.6) in some time interval $[0, T]$ with*

$$V \in C^0([0, T], W^{s,2}(\mathbb{R})) \cap C^1([0, T], W^{s-1,2}(\mathbb{R})),$$

and $T = T(\|V_0\|_{s,2}) > 0$.

Global existence for small data by Tarabek, solution with bounded norm $\|V(t)\|_{s,2}$.

Rewrite the differential equations:

$$u_{tt} - \alpha u_{xx} + \beta \theta_x = h_1 \tag{2.7}$$

$$\theta_t + \gamma q_x + \delta u_{tx} = h_2 \tag{2.8}$$

$$\tau_0 q_t + q + \kappa \theta_x = h_3 \tag{2.9}$$

where

$$\alpha := S_{u_x}(0, 0, 0), \quad \beta := -S_\theta(0, 0, 0), \quad \gamma := a_4(0, 0, 0) \quad (2.10)$$

$$\delta := a_3(0, 0, 0), \quad \kappa := \tau_0 a_6(0, 0, 0), \quad \tau_0 := \tau(0, 0) \quad (2.11)$$

$$h_1 := (S(u_x, \theta, q) - \alpha u_x + \beta \theta)_x \quad (2.12)$$

$$h_2 := (\gamma - a_4)q_x + (\delta - a_3)u_{tx} - \sigma_2 \theta_x - a_5 q \quad (2.13)$$

$$h_3 := \tau_0 \left\{ \left(\frac{1}{\tau_0} - \frac{1}{\tau} \right) q + \left(\frac{\kappa}{\tau_0} - \frac{k}{\tau} \right) \theta_x \right\}. \quad (2.14)$$

One knows

$$\sigma_1(0, 0, 0) = 0, \quad \sigma_2(0, 0, 0) = 0, \quad a_5(0, 0, 0) = 0. \quad (2.15)$$

Defining

$$V := (\sqrt{\alpha \kappa \delta} u_x, u_t, \theta, q)'$$

we obtain

$$A^0 V_t + A^1 V_x + BV = \tilde{F}(V, V_x), \quad V(0) = V_0 := (\sqrt{\alpha\kappa\delta} u_{0,x}, u_1, \theta_0, q_0)' \quad (2.16)$$

with a nonlinearity that vanishes quadratically near zero. The linearized system, i.e. for $\tilde{F} = 0$, is solved by

$$V(t) = e^{tR} V_0$$

where

$$R := -(A^0)^{-1} (A^1 \partial_x + B)$$

generates a C_0 -semigroup on $D(R) := (W^{1,2}(\mathbb{R}))^4 \subset (L^2(\mathbb{R}))^4$. Then the solution to (2.16) is represented in general by

$$V(t) = e^{tR} V_0 + \int_0^t e^{(t-r)R} F(V, V_x)(r) dr \quad (2.17)$$

where

$$F := (A^0)^{-1} \tilde{F} \quad (2.18)$$

Cattaneo's law is assumed to be the linear one:

$$\tau_0 q_x + q + \kappa \theta_x = 0 \quad (2.19)$$

This assumption implies that in (2.9)

$$h_3 = 0 \quad (2.20)$$

Further assumption:

$$\psi(u_x, \theta, q) = \psi_0(u_x, \theta) + \psi_1(u_x, \theta) q^2 \quad (2.21)$$

$$e(u_x, \theta, q) = e_0(u_x, \theta) + e_1(u_x, \theta) q^2 \quad (2.22)$$

Theorem 2.2. *Let $V_0 \in W^{3,2}(\mathbb{R}) \cap L^1(\mathbb{R})$, and $m_0 := \|V_0\|_{3,2} + \|V_0\|_1$. For sufficiently small m_0 there is $c > 0$ such that for $t \geq 0$, the solution V to (2.16) satisfies*

$$\|V(t)\|_{1,2} \leq c(1+t)^{-\frac{1}{4}} \cdot m_0$$

The fact that F vanishes quadratically in zero is sufficient despite the linear decay behavior like that of a heat equation, because terms like " $V \cdot V_x$ " appear. Therefore, the proof rewrites the nonlinearities in divergence form in order to be able to exploit the better decay of derivatives in the later estimates for $\|V(t)\|_2$.

3 Propagation of singularities in three dimensions

(Joint work with Y.-G. Wang)

$$U_{tt}^{po} - \alpha^2 \Delta U^{po} p + \beta \nabla \theta = 0, \quad (3.1)$$

$$\theta_t + \gamma \nabla' q^{po} + \delta \nabla' U_t^{po} = 0, \quad (3.2)$$

$$\tau q_t^{po} + q^{po} + \kappa \nabla \theta = 0, \quad (3.3)$$

$$U^{po}(0, \cdot) = U_0^{po}, \quad U_t^{po}(0, \cdot) = U_1^{po}, \quad \theta(0, \cdot) = \theta^0, \quad q^{po}(0, \cdot) = q_0^{po}. \quad (3.4)$$

The data $(\nabla U^0, U^1, \theta^0, q^0)$ are assumed to be smooth away from and possibly having jumps on a given smooth surface

$$\sigma = \{x \in \Omega_0 \subset \mathbb{R}^3 \mid \Phi^0(x) = 0\}$$

described by a C^2 -function $\Phi^0 : \Omega_0 \subset \mathbb{R}^3 \rightarrow \mathbb{R}$, with Ω_0 open, satisfying

$$|\nabla\Phi^0(x)| = 1 \quad \text{on} \quad \sigma. \quad (3.5)$$

$$\tilde{U} = (\tilde{U}_1, \tilde{U}_2, \tilde{U}_3, \tilde{U}_4)' := (\nabla'U^{po}, U_t^{po}, \theta, q^{po})'.$$

Then

$$\partial_t \tilde{U} + \sum_{j=1}^3 A_j \partial_j \tilde{U} + A_0 \tilde{U} = 0, \quad \tilde{U}(0, \cdot) = \tilde{U}^0, \quad (3.6)$$

The characteristic polynomial

$$\det(\lambda \text{Id}_{\mathbb{R}^8} - \sum_{j=1}^3 \xi_j A_j) = \lambda^4 \left(\lambda^4 - \lambda^2 \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right) |\xi|^2 + \frac{\kappa\gamma\alpha^2}{\tau} |\xi|^4 \right) \quad (3.7)$$

has real roots $\lambda_k, 1 \leq k \leq 8$, and the eigenvalues of

$$B := \sum_{j=1}^3 \partial_j \Phi^0 A_j \quad (3.8)$$

are

$$\lambda_k = 0, \quad 1 \leq k \leq 4, \quad (3.9)$$

$$\lambda_k = \mu_k = \pm \left\{ \frac{1}{2} \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right) \pm \frac{1}{2} \sqrt{ \left(\alpha^2 + \beta\delta + \frac{\kappa\gamma}{\tau} \right)^2 - \frac{4\kappa\gamma\alpha^2}{\tau} } \right\}^{\frac{1}{2}}, \quad 5 \leq k \leq 8 \quad (3.10)$$

taking

$$\mu_8 < \mu_6 < 0 < \mu_7 < \mu_5. \quad (3.11)$$

The characteristic surfaces $\Sigma_k = \{(t, x) \mid \Phi_k(t, x) = 0\}$, $1 \leq k \leq 8$, evolving from the initial surface $\sigma = \{(0, x) \mid \Phi^0(x) = 0\}$ are determined by

$$\partial_t \Phi_k + \mu_k |\nabla \Phi_k| = 0, \quad \Phi_k(0, \cdot) = \Phi^0 \quad (3.12)$$

hence

$$\Sigma_k = \{(t, x) \mid -\mu_k t + \Phi^0(x) = 0\}. \quad (3.13)$$

The associated right (column) eigenvectors r_k and left (row) eigenvectors l_k of B , can be computed explicitly, as well as expansions of these and the eigenvalues in powers of τ . Let the matrices L and R be given by

$$L := \begin{pmatrix} l_1 \\ \vdots \\ l_8 \end{pmatrix}, \quad R := (r_1, \dots, r_8).$$

Then $V := L\tilde{U}$, with \tilde{U} satisfying (3.6), satisfies

$$\partial_t V + \sum_{j=1}^3 (LA_j R) \partial_j V + \tilde{A}_0 V = 0, \quad V(t=0) = V^0 := L\tilde{U}^0, \quad (3.14)$$

where

$$\tilde{A}_0 := LA_0 R + \sum_{j=1}^3 LA_j \partial_j R. \quad (3.15)$$

Let $[H]_{\Sigma_k}$ denote the jump of H along Σ_k . Then (V_1, V_2, V_3, V_4) are continuous at $\cup_{k=5}^8 \Sigma_k$, and $V_j, j = 5, \dots, 8$, does not have any jump on $\Sigma_k, k = 1, \dots, 8$ for $k \neq j$. Moreover, $[V_k]_{\Sigma_0} = 0$ for $1 \leq k \leq 4$. Evolutional equations for $[V_k]_{\Sigma_k}$ for $5 \leq k \leq 8$:

$$\partial_t V_k + \sum_{m=1}^8 \sum_{j=1}^3 (LA_j R)_{km} \partial_j V_m = - \sum_{m=1}^8 (\tilde{A}_0)_{km} V_m + \tilde{F}_k \quad (3.16)$$

for $5 \leq k \leq 8$. Since

$$(\lambda_k \text{Id}_{\mathbb{R}^8} - \sum_{j=1}^3 L \partial_j \Phi^0 A_j R)_{kk} = (\lambda_k \text{Id}_{\mathbb{R}^8} - \tilde{\Lambda})_{kk} = 0 \quad (3.17)$$

with $\tilde{\Lambda} = \text{diag}(\lambda_1, \dots, \lambda_8)$, we obtain that for each $5 \leq k \leq 8$, $[V_k]_{\Sigma_k}$ satisfies the

following transport equation:

$$\left(\partial_t + \sum_{j=1}^3 (LA_j R)_{kk} \partial_{x_j} + (\tilde{A}_0)_{kk}\right) [V_k]_{\Sigma_k} = [\tilde{F}_k]_{\Sigma_k} \quad (3.18)$$

with initial conditions

$$[V_k]_{\Sigma_k}(t = 0) = [V_k^0]_{\sigma}. \quad (3.19)$$

In order to determine the behavior of $[V_k]_{\Sigma_k}$ from (3.18), (3.19), it is essential to study $(\tilde{A}_0)_{kk}$, where, by (3.15), \tilde{A}_0 is given by

$$\tilde{A}_0 = LA_0 R + \sum_{j=1}^3 LA_j \partial_j R. \quad (3.20)$$

$$(\tilde{A}_0)_{kk} = \begin{cases} \frac{1}{2\tau} \pm \frac{1}{2} \sqrt{\frac{\kappa\gamma}{\tau}} \Delta\Phi^0 + \mathcal{O}(1), & k = 5, 6, \\ \frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2} \Delta\Phi^0 + \mathcal{O}(\tau), & k = 7, 8. \end{cases} \quad (3.21)$$

Notice: The mean curvature H of the initial surface σ equals

$$H = \frac{\Delta\Phi^0}{2}$$

and will play an essential role in the behavior of the jumps as $t \rightarrow \infty$ or as $\tau \rightarrow 0$.

$$\begin{aligned} [V_k]_{\Sigma_k(t)} &= [V_k^0]_{\sigma} e^{-\int_0^t (\tilde{A}_0)_{kk}(x(s;0,x^0)) ds} \\ &= \begin{cases} e^{-\frac{t}{2\tau} \mp \frac{1}{2} \sqrt{\frac{\kappa\gamma}{\tau}} \int_0^t (\Delta\Phi^0(x(s;0,x^0)) + O(\sqrt{\tau})) ds} [V_k^0]_{\sigma}, & k = 5, 6, \\ e^{-\frac{\beta\delta}{2\kappa\gamma} t \mp \frac{\alpha}{2} \int_0^t (\Delta\Phi^0(x(s;0,x^0)) + O(\tau)) ds} [V_k^0]_{\sigma}, & k = 7, 8. \end{cases} \end{aligned} \quad (3.22)$$

For $\tau \rightarrow 0$ the dominating term for $k = 5, 6$ is $e^{-\frac{t}{2\tau}}$, i.e., we have exponential decay of the jumps of V_k on Σ_k as $\tau \rightarrow 0$ or $t \rightarrow \infty$ for a fixed small $\tau > 0$. If $k = 7, 8$ the dominating term, for $\tau \rightarrow 0$, is $\exp(-\int_0^t (\frac{\beta\delta}{2\kappa\gamma} \pm \frac{\alpha}{2} \Delta\Phi^0(x(s;0,x^0))) ds)$, whether the jumps of V_k on Σ_k decay exponentially depends on the size of the mean curvature ($= \Delta\Phi^0/2$).

Example: Let σ be the sphere of radius r :

$$\sigma = \{x \in \mathbb{R}^3 \mid |x| = r\} = \{x \mid \Phi^0(x) \equiv r - |x| = 0\}.$$

Then, we have

$$\Sigma_k = \{(t, x) \mid \mu_k t = r - |x|\} = \{(t, x) \mid |x| = r - \mu_k t\}.$$

Spreading surfaces, as $t \rightarrow \infty$, are Σ_6, Σ_8 , and

$$\Delta\Phi^0(x_0) = \frac{2}{|x_0|} = \frac{2}{r} > 0.$$

Thus, as $t \rightarrow +\infty$, $[V_6]_{\Sigma_6}$ is decaying exponentially, while $[V_8]_{\Sigma_8}$ decays (grows resp.) exponentially if

$$\frac{\beta\delta}{\alpha\kappa\gamma} > \Delta\Phi^0 = \frac{2}{r} \quad \left(\frac{\beta\delta}{\alpha\kappa\gamma} < \Delta\Phi^0 = \frac{2}{r} \text{ resp.} \right), \quad (3.23)$$

that is depending on the size of the mean curvature $H = \frac{1}{r}$.

Theorem 3.1. *Suppose that the initial data $\nabla'U^{p_0}, \partial_t U^{p_0}, \theta$ and q^{p_0} may have jumps on $\sigma = \{\Phi^0(x) = 0\}$ with $|\nabla\Phi^0(x)| = 1$, then the propagation of strong singularities of solutions to the linearized problem (3.1)–(3.4) obeys*

(1) *The jumps of $\nabla'U^{p_0}, \partial_t U^{p_0}, \theta, q^{p_0}$ on Σ_5 and Σ_6 decay exponentially both when $\tau \rightarrow 0$ for a fixed $t > 0$ and when $t \rightarrow +\infty$ for a fixed $\tau > 0$.*

(2) *The jumps of $\nabla'U^{p_0}, \partial_t U^{p_0}, q^{p_0}$ on Σ_7 (Σ_8 resp.) are propagated, and when $t \rightarrow +\infty$ they will decay exponentially as soon as $\frac{\beta\delta}{\kappa\gamma} + \alpha\Delta\Phi^0$ ($\frac{\beta\delta}{\kappa\gamma} - \alpha\Delta\Phi^0$ resp.) being positive, more rapidly for smaller heat conductive coefficient $\kappa\gamma$, while the jump of the temperature θ on Σ_7 and Σ_8 vanishes of order $O(\tau)$ when $\tau \rightarrow 0$, which shows a smoothing effect in the system (3.1)–(3.3) when the thermoelastic model with second sound converges to the hyperbolic-parabolic type of classical thermoelasticity.*

4 Low frequency expansion in exterior domains

(Joint work with Y. Naito, Y. Shibata)

Let Ω be an exterior domain in \mathbb{R}^3 with $C^{1,1}$ boundary Γ .

$$\begin{aligned}u_{tt} - \mu\Delta u - (\mu + \lambda)\nabla\operatorname{div} u + \beta\nabla\theta &= 0 \\ \theta_t + \gamma\operatorname{div} q + \delta\operatorname{div} u_t &= 0 \\ \tau_0 q_t + q + \kappa\nabla\theta &= 0\end{aligned}\tag{4.1}$$

in $\Omega \times (0, \infty)$ subject to the initial conditions

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad \theta(x, 0) = \theta_0(x), \quad q(x, 0) = q_0(x) \quad \text{in } \Omega$$

and boundary conditions

$$u = 0, \quad \theta = 0 \quad \text{on } \Gamma \times (0, \infty)$$

Aim: Low frequency expansion of the corresponding resolvent problem(s) — important to investigate the decay property of solutions to (4.1) as time goes to infinity (essentially: via Laplace transform).

Moreover: Limit as τ_0 tends to zero.

Resolvent problem:

$$\begin{aligned}
 k^2 u - \mu \Delta u - (\mu + \lambda) \nabla \operatorname{div} u + \beta \nabla \theta &= f && \text{in } \Omega \\
 k\theta + \gamma \operatorname{div} q + \delta k \operatorname{div} u &= g && \text{in } \Omega \\
 \tau_0 k q + q + \kappa \nabla \theta &= h && \text{in } \Omega \\
 u = 0, \quad \theta = 0 &&& \text{on } \Gamma
 \end{aligned} \tag{4.2}$$

First $\Omega = \mathbb{R}^3$: Use Fourier transform. Eliminate q and consider

$$\begin{aligned}
k^2 u - \mu \Delta u - (\mu + \lambda) \nabla \operatorname{div} u + \beta \nabla \theta &= f & \text{in } \Omega \\
k \theta - \gamma \kappa (\tau_0 k + 1)^{-1} \Delta \theta + \delta k \operatorname{div} u &= g & \text{in } \Omega \\
u = \theta = 0 && \text{on } \Gamma
\end{aligned} \tag{4.3}$$

Theorem 4.1. *Let $1 < q < \infty$ and $0 < \tau_0 \leq 1$. Then, for any small $\epsilon > 0$ there exist a constant $\sigma_0 > 0$ depending on ϵ and an operator $S_k \in \operatorname{Anal}(U_{\sigma_0, \epsilon}, \mathcal{B}(L_q(\mathbb{R}^3)^3 \times L_q(\mathbb{R}^3), W_q^2(\mathbb{R}^3)^3 \times W_q^2(\mathbb{R}^3)))$ such that for any $(f, g) \in L_q(\mathbb{R}^3)^3 \times L_q(\mathbb{R}^3)$, $(u, \theta) = S_k(f, g)$ solves equation (4.3). Here, for two Banach spaces X and Y , $\mathcal{B}(X, Y)$ denotes the set of all bounded linear operators from X into Y , $U_{\sigma_0, \epsilon}$ denotes an open set in \mathbb{C} defined by the formula:*

$$U_{\sigma_0, \epsilon} = \{k \in \mathbb{C} \setminus \{0\} \mid |\arg k| \leq (\pi/2) - \epsilon, \quad |k| \leq \sigma_0\}$$

and $\text{Anal}(U_{\sigma_0, \epsilon}, X)$ denotes the set of all holomorphic functions defined on $U_{\sigma_0, \epsilon}$ with their values in X .

Theorem 4.2. *Let $1 < q < \infty$, $0 < \epsilon < \pi/2$, $0 < \tau_0 \leq 1$ and $R > 0$. Let σ_0 and S_k be the same number and solution operator as in Theorem 4.1, respectively. Set*

$$\begin{aligned}\mathcal{L}_{q,R}(\mathbb{R}^3) &= \{(f, g) \in L_q(\mathbb{R}^3)^3 \times L_q(\mathbb{R}^3) \mid (f, g) \text{ vanishes for } |x| > R\} \\ \mathcal{W}_{q,\text{loc}}(\mathbb{R}^3) &= W_{q,\text{loc}}^2(\mathbb{R}^3)^3 \times W_{q,\text{loc}}^2(\mathbb{R}^3)\end{aligned}$$

Then, there exist a σ ($0 < \sigma \leq \sigma_0$) and $G_j(k) \in \text{Anal}(U_\sigma, \mathcal{B}(\mathcal{L}_{q,R}(\mathbb{R}^3), \mathcal{W}_{q,\text{loc}}(\mathbb{R}^3)))$ ($j = 0, 1$) such that when $(f, g) \in \mathcal{L}_{q,R}(\mathbb{R}^3)$, $G_k(f, g) = (k^{1/2}G_0(k) + G_1(k))(f, g)$ solves equation (4.3) for $k \in U_\sigma$ and $G_k(f, g) = S_k(f, g)$ for $k \in U_{\sigma, \epsilon}$.

Now let Ω be an exterior domain.

Let σ , S_k , $G_0(k)$ and $G_1(k)$ be the same constant and operators as in Theorem 4.2

and set

$$G_k = k^{1/2}G_0(k) + G_1(k) \quad (4.4)$$

Locally we get

Theorem 4.3. *Let $1 < q < \infty$ and $0 < \tau_0 \leq 1$. Let R be a large fixed number such that $\mathbb{R}^3 \setminus \Omega \subset B_R$. Then, there exists a small number σ' ($0 < \sigma' \leq \sigma$) and an operator $H_k \in \mathcal{B}(\mathcal{L}_{q,R}, \mathcal{W}_{q,\text{loc}}^2(\Omega))$ for each $k \in U_{\sigma'} = \{k \in \mathbb{C} \mid |k| \leq \sigma'\}$ such that $H_k(f, g)$ satisfies equation (4.3) for any $(f, g) \in \mathcal{L}_{q,R}(\Omega)$ and $k \in U_{\sigma'}$ and H_k has the expansion formula:*

$$H_k = k^{1/2}H_0(k) + H_1(k) \quad \text{for } k \in U_{\sigma'}$$

where $H_k^0, H_k^1 \in \text{Anal}(U_{\sigma'}, \mathcal{B}(\mathcal{L}_{q,R}, \mathcal{W}_{q,\text{loc}}^2(\Omega)))$.

Combining this with the result for $\Omega = \mathbb{R}^3$ by cut-off techniques yields

Theorem 4.4. *Let $1 < q < \infty$, $0 < \epsilon < \pi/2$ and $0 < \tau_0 \leq 1$. Let $\sigma' > 0$ be the same constant as in Theorem 4.3. Then, there exists an operator $T_k \in \text{Anal}(U_{\sigma', \epsilon}, \mathcal{B}(L_q(\Omega)^4, W_q^2(\Omega)^4))$ such that $T_k(f, g)$ satisfies equation (4.3) for any $(f, g) \in L_q(\Omega)^4$ and $k \in U_{\sigma', \epsilon}$.*

Employing the same argument, we can show the theorems corresponding to Theorems 4.3 and 4.4 in the classical thermoelastic case ($\tau_0 = 0$). Moreover, we have

Theorem 4.5. *The solution operators H_k constructed in Theorem 4.3 and T_k in Theorem 4.4 depend on $\tau_0 \in (0, 1]$ continuously. The limit of H_k and T_k as $\tau_0 \rightarrow 0$ are the corresponding operators of the classical thermoelastic equations, where the limit is given in the operator norm of $\mathcal{B}(\mathcal{L}_{q,R}(\Omega), \mathcal{W}_{q,\text{loc}}^2(\Omega))$ when $k \in U_{\sigma'}$ and $\mathcal{B}(L_q(\Omega)^4, W_q^2(\Omega)^4)$ when $\text{Re } k > 0$ and $|k| < \sigma'$, respectively.*

Thank you for your attention!