# On the Cahn-Hilliard equation with singular potential and dynamic boundary conditions

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#### DICOP '08 – Direct, Inverse and Control Problems Cortona, September 24, 2008

Introduction and motivation

Well-posedness

Long-time behavior

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We deal with the following equation of Cahn-Hilliard type:

$$u_t - \Delta (-\Delta u + f(u) - h) = 0,$$
 (CH)

where

- the unknown u represents an order parameter,
- ►  $w = -\Delta u + f(u) h$  is the chemical potential, and we assume  $\partial_n w = 0$ ,
- ▶ *f* is the derivative of a configuration potential *F*,
- h is an external source, included in view of possible applications to conserved phase field systems.

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- Usually, equation (CH) is coupled with either homogeneous Neumann or periodic boundary condition.
- Recently, the following class of dynamic boundary conditions has been proposed:

$$v_t - \Delta_{\Gamma} v + f_{\Gamma}(v) = h_{\Gamma} - \partial_n u$$
, where (dynBC)

- v represents the trace of u on  $\Gamma = \partial \Omega$ ,
- > ∆<sub>r</sub> is the Laplace-Beltrami operator
- $b_{\rm F}$  is a boundary source,
- F is the derivative of a boundary configuration potential F-
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- In a variational derivation, it comes from a boundary contribution to the free energy;
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- The potentials F,  $F_{\Gamma}$  are assumed to be  $\lambda$ -convex;
- Two situations appear to be physically relevant:
  - Regular potentials: when F (or Fr) is a smooth function defined on the whole real line (example: F(r) ~ (r<sup>2</sup> = 1)<sup>2</sup>);
     Singular potentials: when F (or Fr) has a bounded domain (e.g., (-1,1) or (-1,1)) and it is conventionally set as 4.5, outside it (example: F(r) ~ (1 = r)) (or (1 = r)).
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- Consider the widest admissible class of potentials, both on the "bulk" Ω and on the boundary Γ;
- Analyze the growth and/or compatibility compatibility conditions which are required in order to have (at least) existence of a solution;
- Prove a well-posedness theorem via a rigorous approximation argument;
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#### Main trouble:

A control of f(u),  $f_{\Gamma}(v)$  cannot be deduced from a bound on u

(the latter holds due to the energy inequality coming from the variational structure of the problem).

But this is true also, e.g., for the Neumann conditions (cf., e.g., [Kenmochi-Niezgodka-Pawłow]);

Is this really a new difficulty?

YES! The coupling between singular potentials and dynamic b.c. creates additional problems.

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## The key point

#### Let us detail the estimate for f(u)



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... but it is better to do it on the blackboard!



#### Although the free energy is of the form

$$E(u, v) = \int_{\Omega} F(u) + \int_{\Gamma} F_{\Gamma}(v) + \dots,$$

in order to control f(u) (and  $f_{\Gamma}(v)$ ) uniformly in time, we need an information on

and not only on  $F_{\Gamma}(v_0)$ .

Thus, the phase space will be (a bit) smaller than the energy space.

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#### We need some condition ensuring that

$$\int_{\Gamma} f(v) (f_{\Gamma}(v) - h_{\Gamma}) \quad \text{can be controlled};$$

- ► either F<sub>Γ</sub> is singular "in a similar way"
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- Then, there exists a unique solution to (CH)+(dynBC)
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- Notice that, due to the Neumann condition for w = −∆u + f(u) − h, the spatial average of u is conserved in time. Let us call m its (known) value;
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$$-\Delta u_{\infty} + f(u_{\infty}) = w_{\infty} + h, \qquad |\Omega|^{-1} \int_{\Omega} u_{\infty} = m, \ (CH_{\infty})$$

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- Note that  $w_{\infty}$  is a constant function.
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- Can we, for analytic f and f<sub>Γ</sub> prove, via the Łojasiewicz-Simon inequality that the ω-limit is a singleton?
- ► For singular potentials *F*, we need that at least the stationary states u<sub>∞</sub> are "separated" from the boundary of dom *f*. Namely, if dom *f* = (-1, 1), we need

 $\exists \varepsilon > 0: -1 + \varepsilon \leq u_{\infty}(x) \leq 1 - \varepsilon, \quad \text{a.e. in } \Omega.$ 

▶ We can prove this for *f* such that (additionally)

 $\lim_{r\to\partial\,\mathrm{dom}\,f}|f(r)|=+\infty$ 

Main ingredient: the control of w<sub>ee</sub>.

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## References

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Thanks for your attention



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