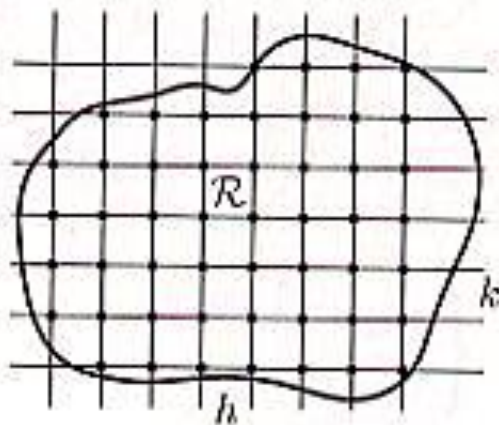


Numerical Methods for Partial Differential Equations (PDE) (2) Finite Difference Methods

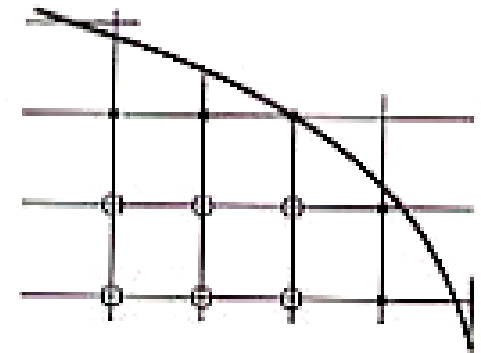
PDE - Finite Differences

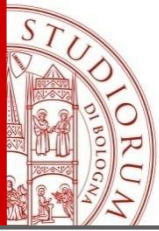
Basic Idea:

- Replace the region R with a (rectangular) grid (mesh) of points of R
- “collocate” the differential equation on the nodes of the mesh
- approximate (at nodes) the partial derivatives with numerical differentiation formulas (finite differences)



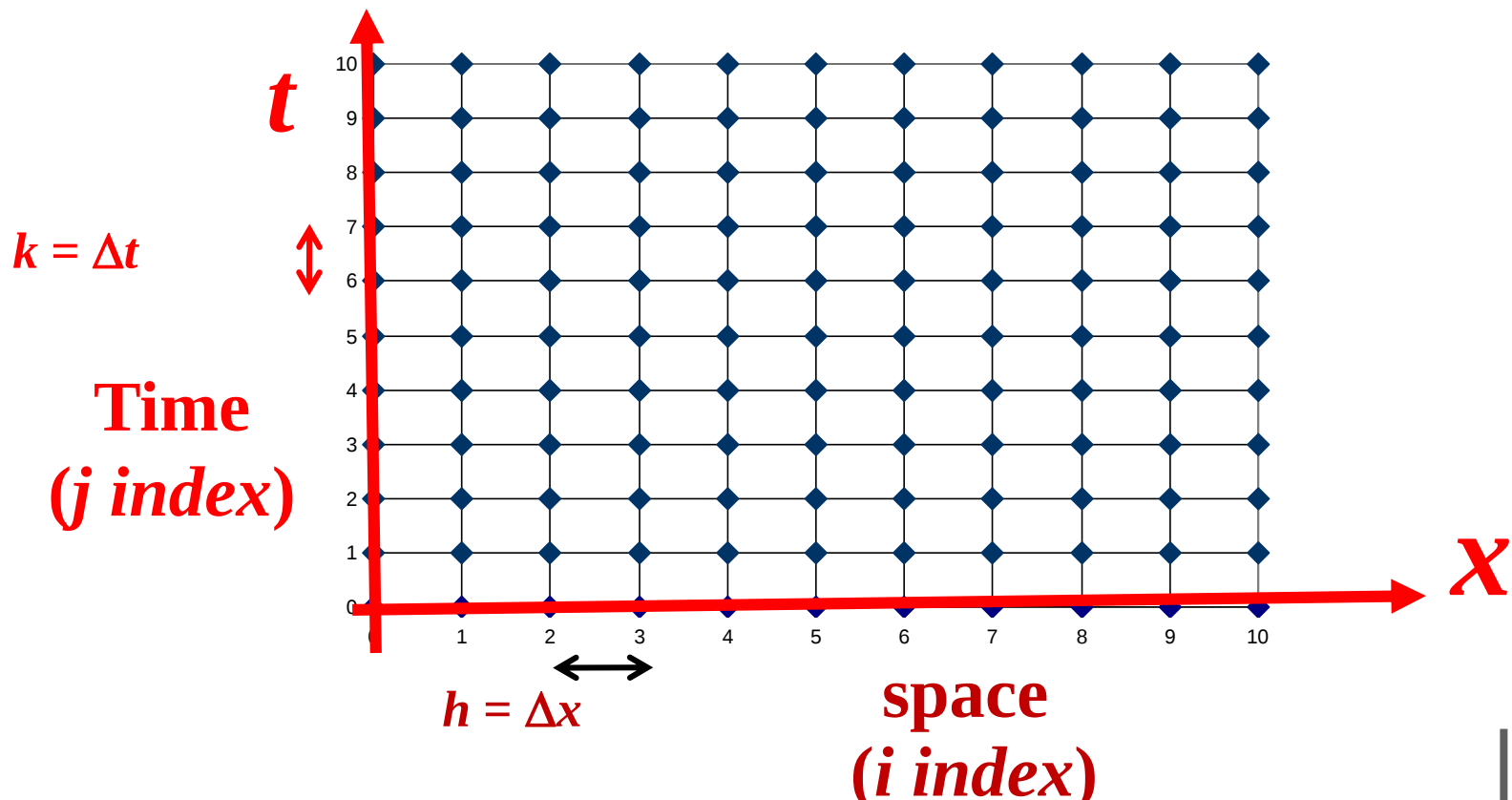
- inner points
- boundary points





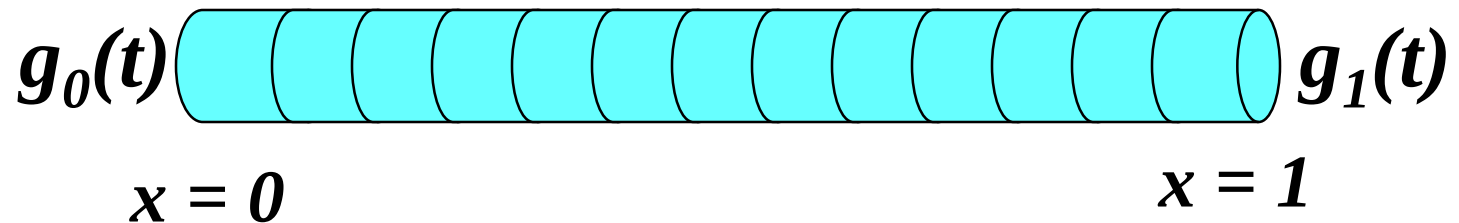
Domain Discretization

- Define the grid of points $(x_i, t_j) = (i\Delta x, j\Delta t)$,
- Discretize the continuous function $u(x, t)$: $u_{i,j} = u(i\Delta x, j\Delta t)$.



1) Parabolic PDE 1D of order 2 heat or diffusion

Wire (of unit length) thermally isolated, the initial distribution of known temperature. The ends of the wire are maintained at known temperatures, in each instant $t > 0$. Compute the temperature distribution $u(x,t)$ at subsequent instants to the initial time.

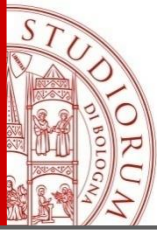


$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}; \quad 0 < x < 1, \quad 0 \leq t \leq T$$

$$\text{CI: } u(x, 0) = f(x), \quad 0 < x < 1$$

$$\text{CB: } \begin{cases} u(0, t) = g_0(t) \\ u(1, t) = g_1(t) \end{cases}, \quad 0 \leq t \leq T$$

Dirichlet



Parabolic PDE : Heat or Diffusion

Well posed problem

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}; \quad 0 \leq t \leq T$$

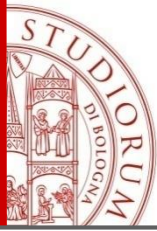
Init. Cond. $u(x, 0) = f(x), \quad -\infty < x < \infty$

Exact Solution

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} e^{-(x-s)^2 / 4t} f(s) ds$$

Smoothing effect: e.g., even if the initial datum $f(s)$ is only bounded and locally piecewise continuous in the x direction, the solution $u(x, t)$ is infinitely derivable for each $t > 0$. Moreover $\lim_{\substack{x \rightarrow \xi \\ t \rightarrow 0}} u(x, t) = f(\xi)$

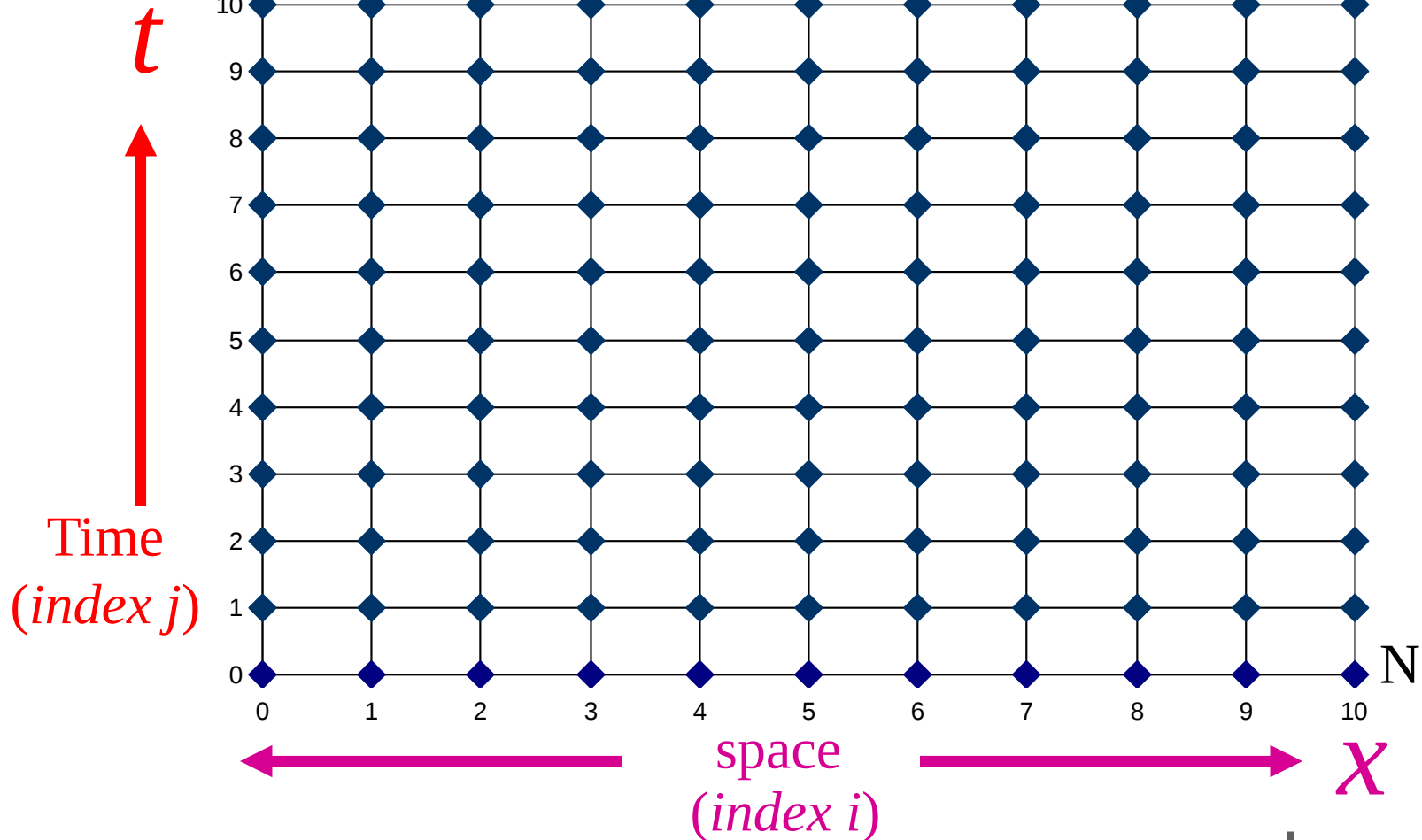
This behaviour is typical of parabolic PDEs.



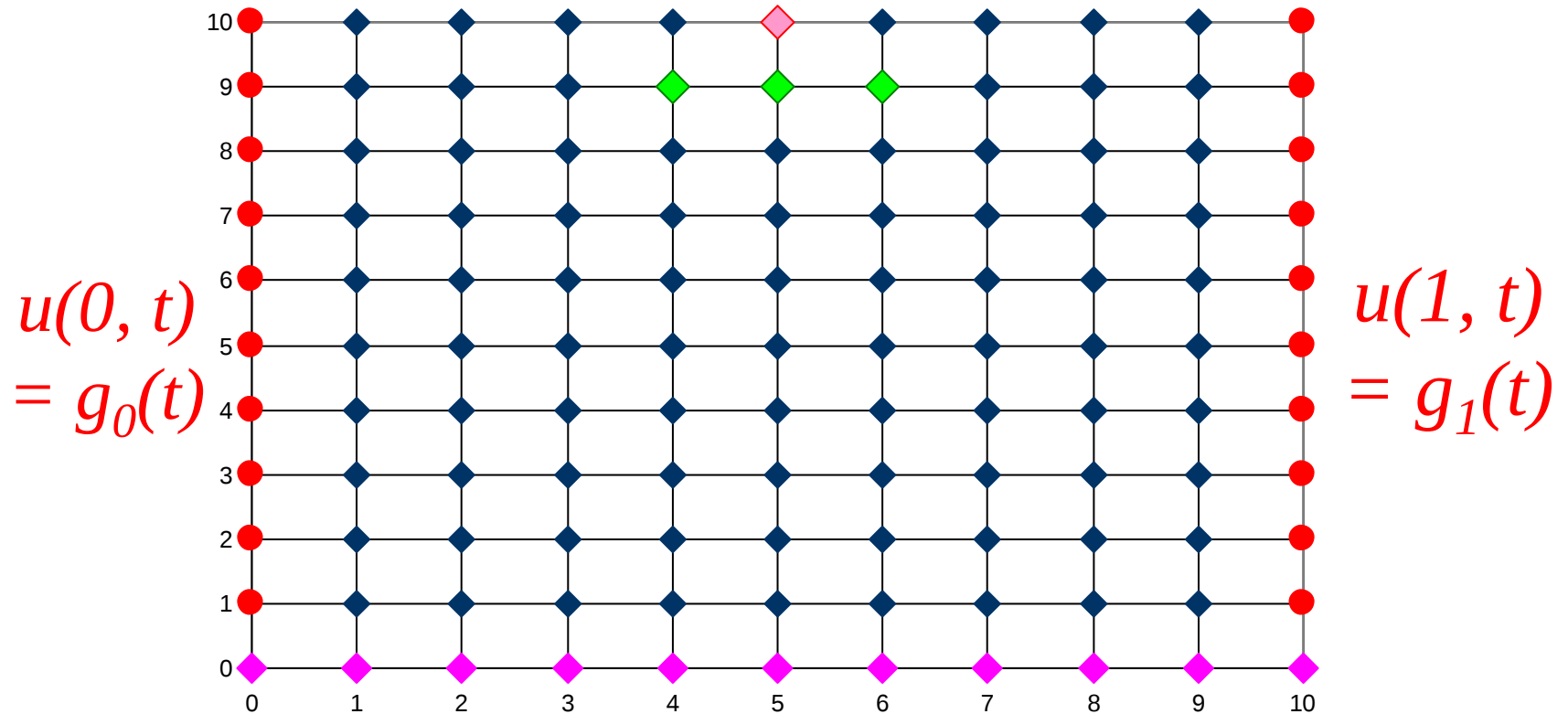
Parabolic PDE : Heat or Diffusion

$$h = \Delta x$$

$$k = \Delta t$$



Boundary and initial conditions

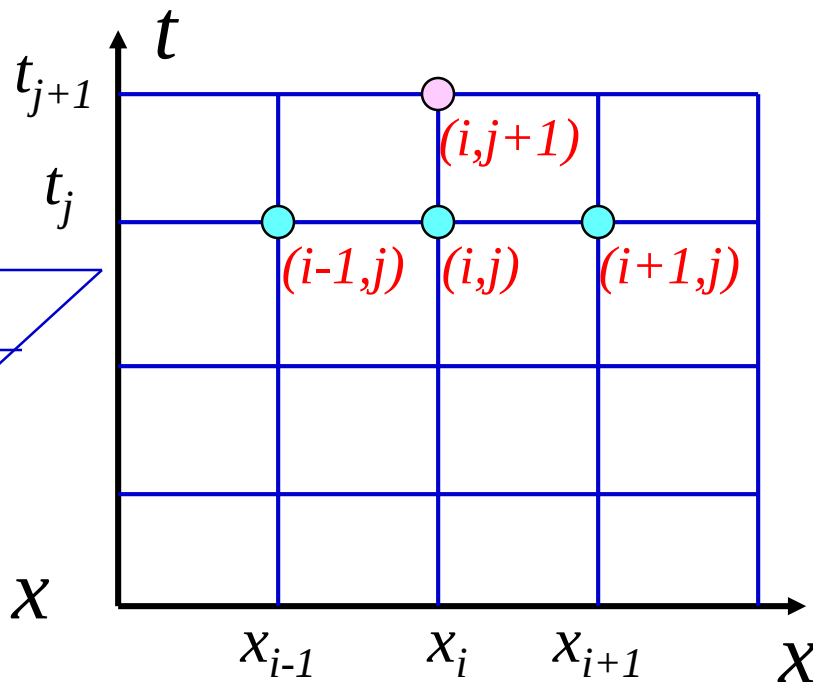
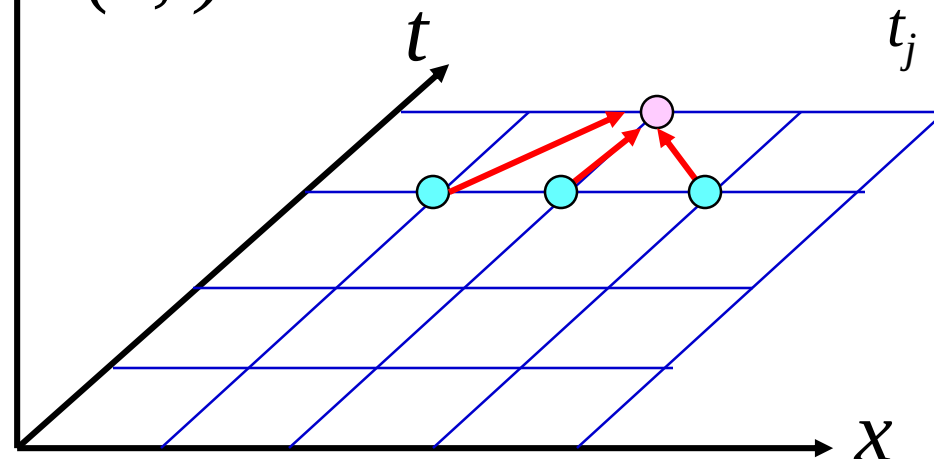


Initial conditions: $u(x, 0) = f(x)$

Explicit method

Let $\begin{cases} h = \Delta x = 1/N, & x_i = ih \\ k = \Delta t = T/M, & t_j = jk \end{cases}$

$u(x, t)$

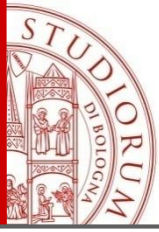


Forward Difference

$$\frac{\partial u}{\partial t} (x_i, t_j) = \frac{1}{k} (u_i^{j+1} - u_i^j) + O(k)$$

Centered Difference
at time j

$$\frac{\partial^2 u}{\partial x^2} (x_i, t_j) = \frac{1}{h^2} (u_{i-1}^j - 2u_i^j + u_{i+1}^j) + O(h^2)$$



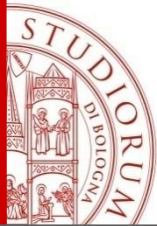
Explicit method

$$u_t = u_{xx}$$

$$\frac{1}{k}(u_{i,j+1} - u_{i,j}) = \frac{1}{h^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j})$$

$$\begin{aligned} u_{i,j+1} &= u_{i,j} + \frac{k}{h^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) \\ &= ru_{i-1,j} + (1 - 2r)u_{i,j} + ru_{i+1,j} \end{aligned}$$

$$r = \frac{k}{h^2} = \frac{\Delta t}{\Delta x^2}$$



$$CI : u_{i,0} = u_0(x_i) \quad i = 0, 1, \dots, n$$

$$CB : u_{0,t} = g_0(t), \quad u_{n,t} = g_1(t)$$

$$\begin{bmatrix} u_{1,j+1} \\ \vdots \\ \vdots \\ u_{N-1,j+1} \end{bmatrix} = \begin{bmatrix} 1-2r & r & & 0 \\ r & 1-2r & r & \\ & \ddots & \ddots & r \\ 0 & & r & 1-2r \end{bmatrix} \begin{bmatrix} u_{1,j} \\ \vdots \\ \vdots \\ u_{N-1,j} \end{bmatrix} + \begin{bmatrix} ru_{0,j} \\ 0 \\ 0 \\ ru_{N,j} \end{bmatrix}$$

$$U_{j+1} = AU_j + v \quad A \in \mathbb{R}^{(N-1) \times (N-1)} \quad v, U \in \mathbb{R}^{(N-1)}$$

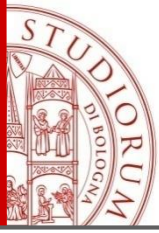
Conditioned Stability:

$$0 < r \leq 0.5$$

$$\begin{bmatrix} u_{1,j+1} \\ \vdots \\ \vdots \\ u_{N-1,j+1} \end{bmatrix} = \left(\begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ & 0 & \ddots & 0 \\ 0 & & 0 & 1 \end{bmatrix} + r \begin{bmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ & \ddots & \ddots & 1 \\ 0 & & 1 & -2 \end{bmatrix} \right) \begin{bmatrix} u_{1,j} \\ \vdots \\ \vdots \\ u_{N-1,j} \end{bmatrix} + \begin{bmatrix} ru_{0,j} \\ 0 \\ \vdots \\ 0 \\ ru_{N,j} \end{bmatrix}$$

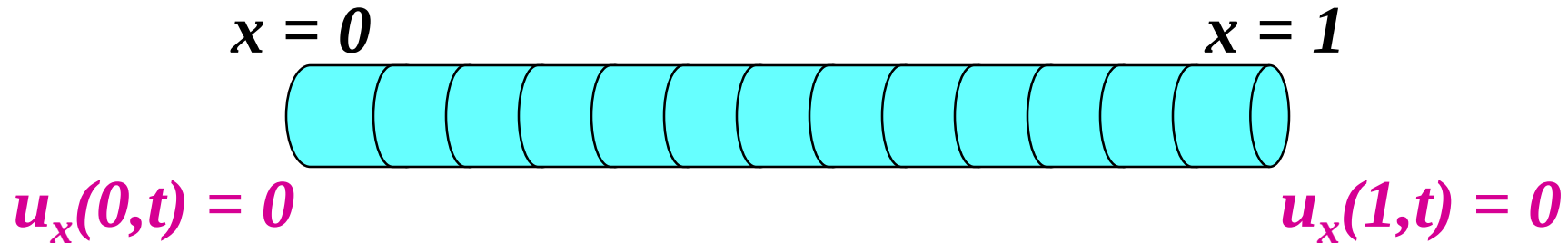
$$U_{j+1} = (I + rL)U_j + v \quad L \in R^{(N-1) \times (N-1)} \quad v, U \in R^{(N-1)}$$

$\frac{1}{h^2} L$ second order derivative matrix, discretization of the Laplace operator



Heat equation and Neumann boundary cond.

No heat flow at boundaries $x = 0$ e $x = 1$

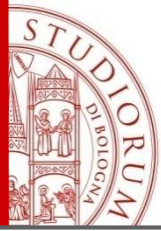


$$\frac{\partial u}{\partial t} = c \frac{\partial^2 u}{\partial x^2}; \quad 0 < x < 1, \quad 0 \leq t \leq T$$

$$u(x,0) = f(x), \quad 0 < x < 1$$

CB:
Neumann

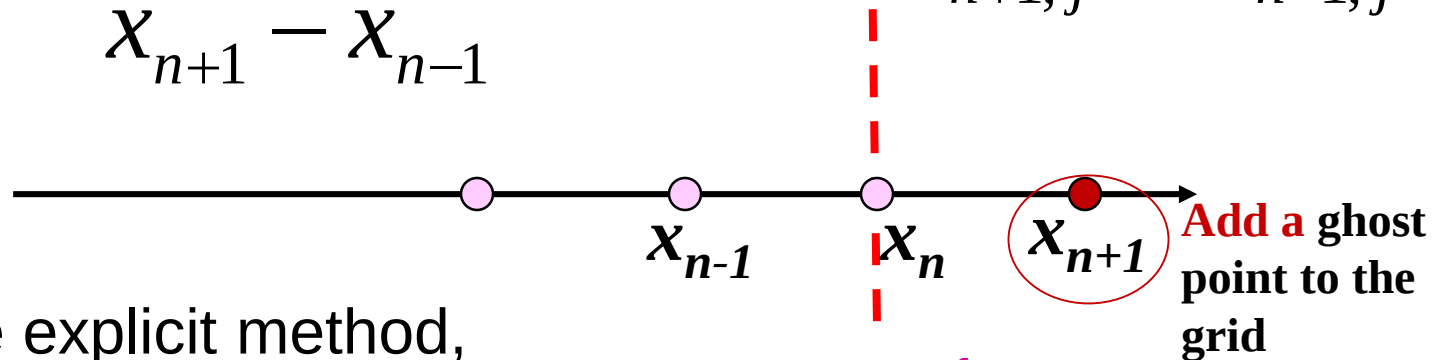
$$\begin{cases} u_x(0,t) = 0 \\ u_x(1,t) = 0 \end{cases}, \quad 0 \leq t \leq T$$



Homogeneous Neumann Boundary Conditions: isolated bounds

No heat flow at boundaries $x = 1$ $u_x(1,t)=0$

$$u_x(1, t_j) = \frac{u_{n+1,j} - u_{n-1,j}}{x_{n+1} - x_{n-1}} = 0 \Rightarrow u_{n+1,j} = u_{n-1,j}$$

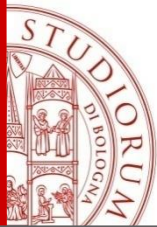


Applying the explicit method,
if $i = n$, we have an extra equation:

$x = 1$

$$\begin{aligned} u_{n,j+1} &= ru_{n-1,j} + (1 - 2r)u_{n,j} + ru_{n+1,j} \\ &= 2ru_{n-1,j} + (1 - 2r)u_{n,j} \end{aligned}$$

$$r = c \frac{k}{h^2} = c \frac{\Delta t}{\Delta x^2}$$



$$CI : u_{i,0} = u_0(x_i) \quad i = 0, 1, \dots, n$$

$$CB : u_x(0, t) = u_x(1, t) = 0$$

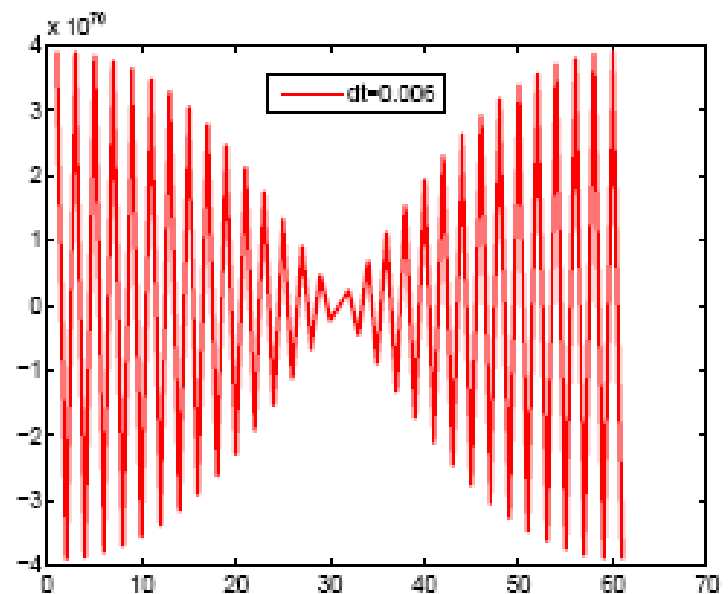
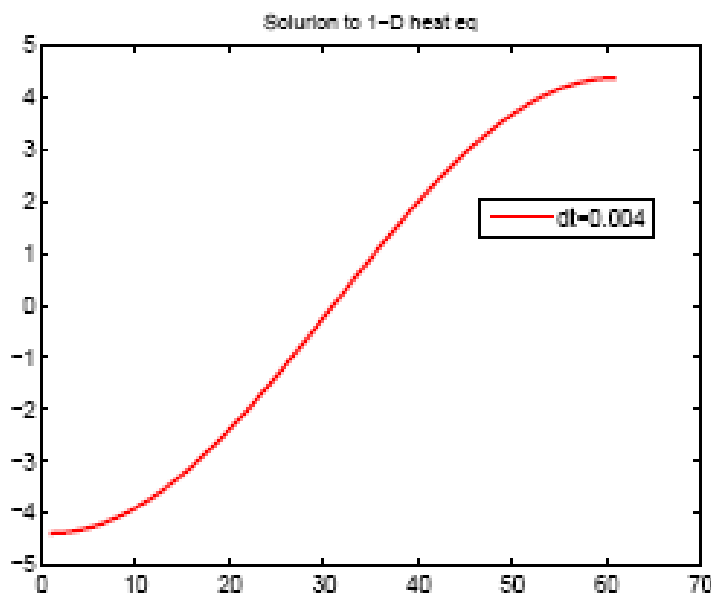
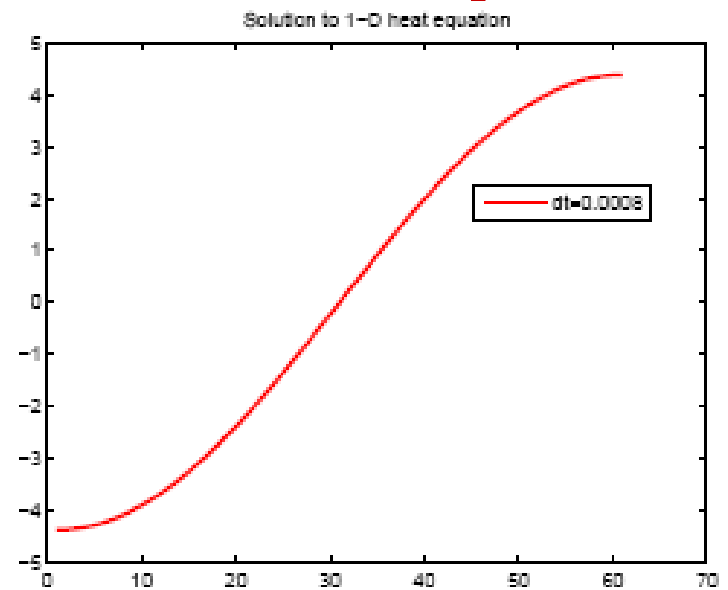
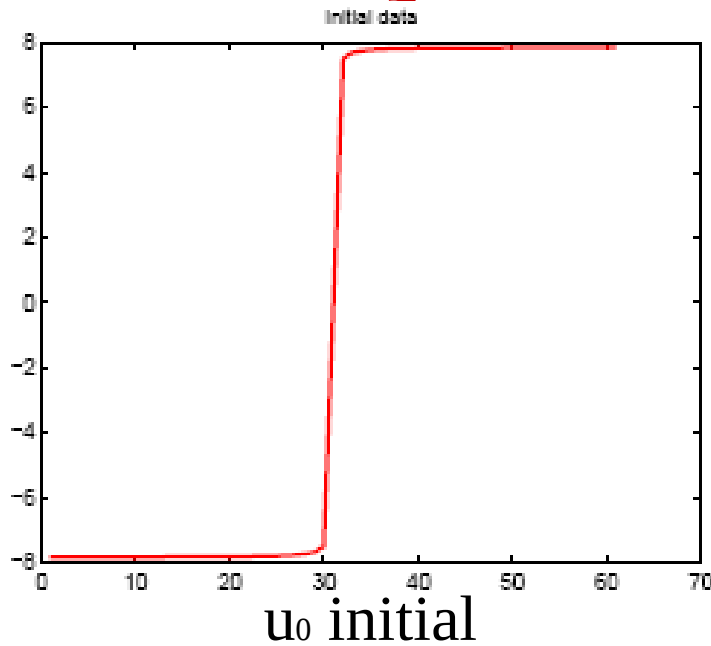
$$\begin{bmatrix} u_{0,j+1} \\ u_{1,j+1} \\ \vdots \\ u_{i,j+1} \\ \vdots \\ u_{N-1,j+1} \\ u_{N,j+1} \end{bmatrix} = \begin{bmatrix} 1-2r & 2r & 0 & & \\ r & 1-2r & r & & \\ & \ddots & \ddots & \ddots & \\ & & 2r & 1-2r & \end{bmatrix} \begin{bmatrix} u_{0,j} \\ u_{1,j} \\ \vdots \\ u_{i,j} \\ \vdots \\ u_{N-1,j} \\ u_{N,j} \end{bmatrix}$$

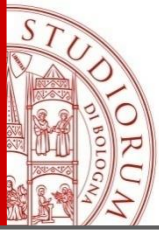
$$U_{j+1} = AU_j \quad A \in \mathbb{R}^{(N+1) \times (N+1)}, U \in \mathbb{R}^{N+1}$$

$$\text{Truncation Error} \quad O(k + h^2)$$

$$\text{Conditioned Stability: } 0 < r \leq 0.5$$

Stability: example heat eq.





Local Truncation Error and order of accuracy

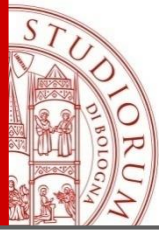
- **Consistency**: For a given numerical scheme, the local truncation error is the error that is generated by demanding that the exact solution verifies the numerical scheme itself.

If the truncation error
$$\tau(k, h) = \max_{i,j} |\tau_i^j| \quad k = \Delta t$$

tends to zero as h and k tend to zero, independently, then the numerical scheme will be defined **consistent**.

- A numerical scheme is said to be **order p accurate in space and order q accurate in time** (for suitable integers p and q), if for a sufficiently smooth solution of the exact problem, we have

$$\tau(k, h) = O(k^p + h^q)$$



Stability and convergence

- A method is said to be **convergent** (in max-norm) if

$$\lim_{k, h \rightarrow 0} (\max |y(x_i, t^j) - u_{i,j}|) = 0$$

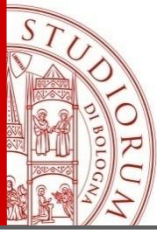
Consistency + stability $\Rightarrow$ convergence

- Stability**

Perturbations on the initial data (initial conditions) are not amplified by the numerical procedure in the solution.

Numerical scheme in general is

$$(*) \quad U_{j+1} = A(k)U_j + b_j(k)$$
$$A(k) \in R^{N \times N} \quad \text{space grid } h = \frac{1}{N+1} \quad b(k) \in R^N$$



Convergence

Def. A linear method of the form (*) is Lax-Richtmyer stable if, for each time T , there is a constant $C_T > 0$ such that

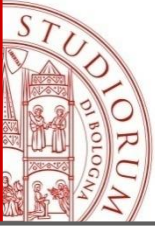
$$\exists C_T > 0 \quad s.t.$$

$$\|A(k)^j\| \leq C_T$$

$$\forall k > 0, \text{ and integers } j \text{ for which } kj \leq T$$

Lax Equivalence Theorem

A consistent linear method of the form (*) is convergent if and only if it is Lax-Richtmyer stable.



Stability of the Explicit Method

Exact solution at time step j : U_j

U_{j+1} is given by multiplying the tridiagonal matrix A for the solution at the previous time step j -th:

$$U_{j+1} = AU_j$$

$$U_{j+1} = A^2 U_{j-1} = \dots = A^{j+1} U_0$$

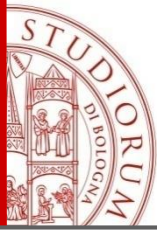
Let us introduce some perturbations E_0 on the data U_0

$$\bar{U}_0 = U_0 + E$$

**Computed
Solution after
 $j+1$ steps**

$$\bar{U}_{j+1} = A^{j+1} \bar{U}_0 = A^{j+1} (U_0 + E) = A^{j+1} U_0 + A^{j+1} E$$

The effect of the error E after $j+1$ steps is $A^{j+1}E$



Stability of the Explicit Method

$$\|A^{j+1}E\| \leq |\lambda|^{j+1} \|E\|$$

$$\|A\|_2 = \rho(A) = \max_{1 \leq p \leq N+1} |\lambda_p|$$

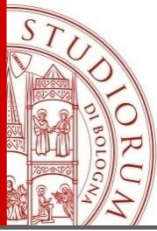
where λ is the dominant eigenvalue of A

- The method is strongly stable when $\|A\|_2 \leq 1$
- Moreover, if the eigenvalues of A are all distinct, the method is stable if and only if

$$\rho(A) \leq 1$$

That is

$$|\lambda| \leq 1$$



Stability of the Explicit Method

The matrix A can be decomposed in the form $A = I + rL$ where

$$L = \begin{bmatrix} -2 & 1 & & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & -2 & 1 \\ & & & 1 & -2 \end{bmatrix}$$

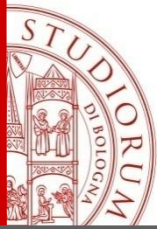
The eigenvalues of L are $\lambda_p = 2(\cos(p\pi h) - 1)$ $p = 1, \dots, N-1, h = 1/N$
 $\lambda_{\max} = -4$

Then, the eigenvalues of A satisfy : $|1 + r\lambda| \leq 1$

The method is then stable when that is $-2 \leq -4 \frac{k}{h^2} \leq 0$

$$r \leq 0.5$$

$$\frac{k}{h^2} \leq \frac{1}{2} \rightarrow k(= \Delta t) \leq \frac{h^2}{2}$$



Example: Explicit Method

$$u_{i,j+1} = ru_{i-1,j} + (1 - 2r)u_{i,j} + ru_{i+1,j}$$

- Stable

$$r = 0.01 \Rightarrow u_{i,j+1} = 0.01 u_{i-1,j} + 0.98 u_{i,j} + 0.01 u_{i+1,j}$$

$$r = 0.1 \Rightarrow u_{i,j+1} = 0.1 u_{i-1,j} + 0.8 u_{i,j} + 0.1 u_{i+1,j}$$

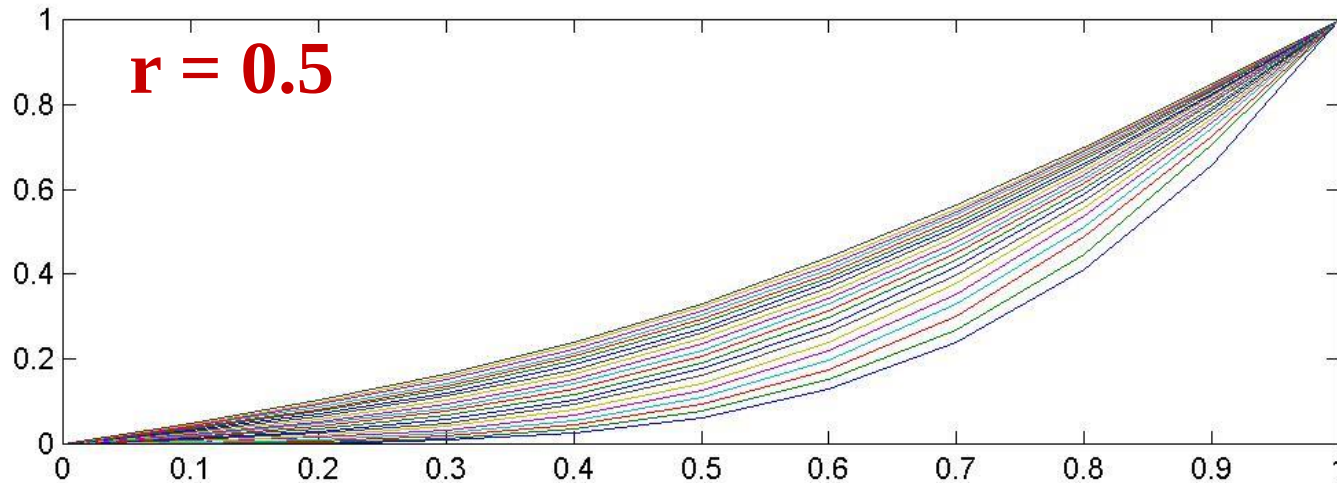
$$r = 0.4 \Rightarrow u_{i,j+1} = 0.4 u_{i-1,j} + 0.2 u_{i,j} + 0.4 u_{i+1,j}$$

$$r = 0.5 \Rightarrow u_{i,j+1} = 0.5 u_{i-1,j} + 0.5 u_{i+1,j}$$

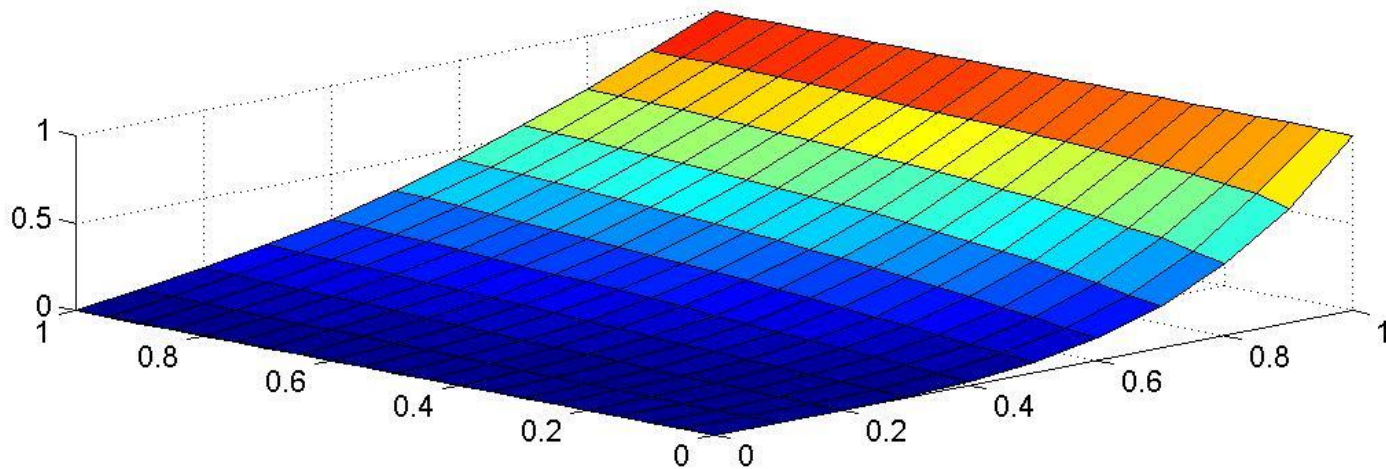
- Unstable (negative coefficients)

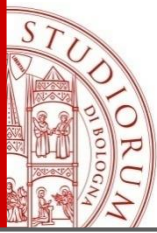
$$\begin{cases} r = 1 & \Rightarrow u_{i,j+1} = u_{i-1,j} - u_{i,j} + u_{i+1,j} \\ r = 10 & \Rightarrow u_{i,j+1} = 10 u_{i-1,j} - 19 u_{i,j} + 10 u_{i+1,j} \\ r = 100 & \Rightarrow u_{i,j+1} = 100 u_{i-1,j} - 199 u_{i,j} + 100 u_{i+1,j} \end{cases}$$

Heat equation: explicit method



Accurato dell'ordine di $O(h^2)$, poichè richiesto $k=O(h^2)$ per stabilità





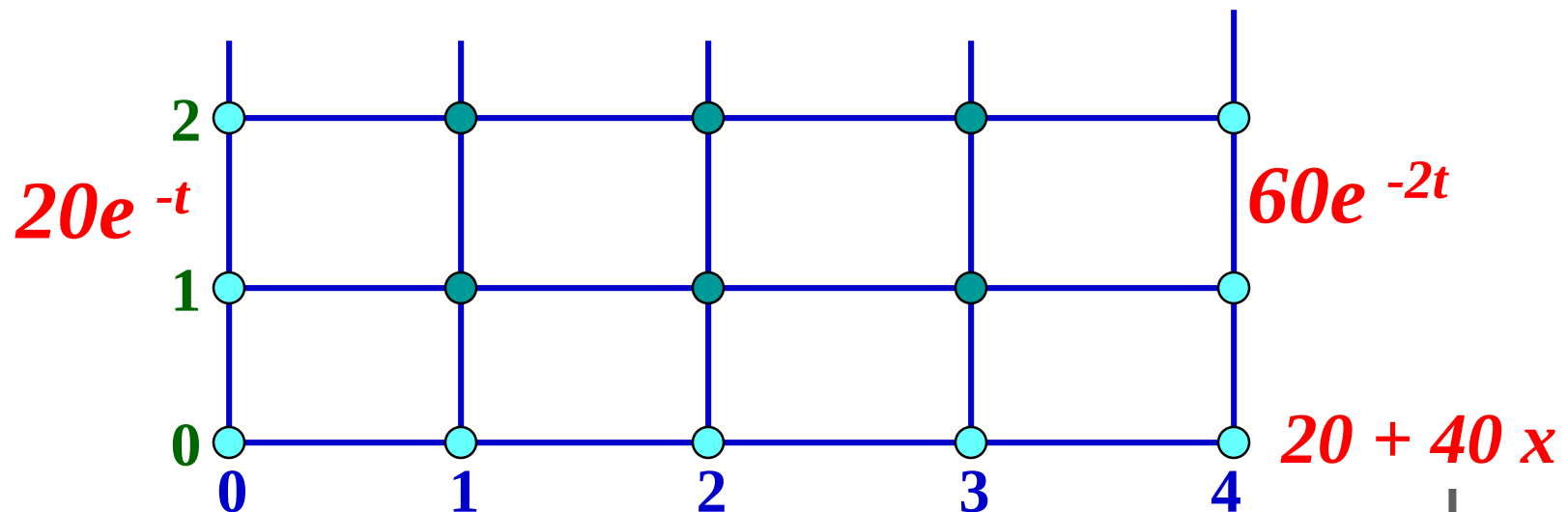
Example: explicit method, stable solution

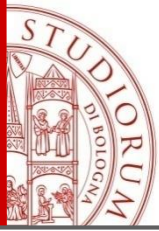
- Heat Equation (Parabolic PDE)

$$u_t = cu_{xx}; \quad 0 \leq x \leq 1$$

$$\begin{cases} u(x,0) = 20 + 40x \\ u(0,t) = 20e^{-t}, \quad u(1,t) = 60e^{-2t} \end{cases}$$

$$c = 0.5, \quad h = 0.25, \quad k = 0.05$$





Example: explicit method, stable solution

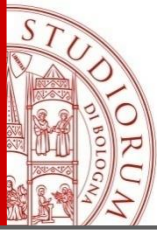
- Compute r

$$r = \frac{ck}{h^2} = \frac{(0.5)(0.05)}{(0.25)^2} = 0.4$$

$$\begin{aligned} u_{i,j+1} &= ru_{i-1,j} + (1 - 2r)u_{i,j} + ru_{i+1,j} \\ &= 0.4u_{i-1,j} + 0.2u_{i,j} + 0.4u_{i+1,j} \end{aligned}$$

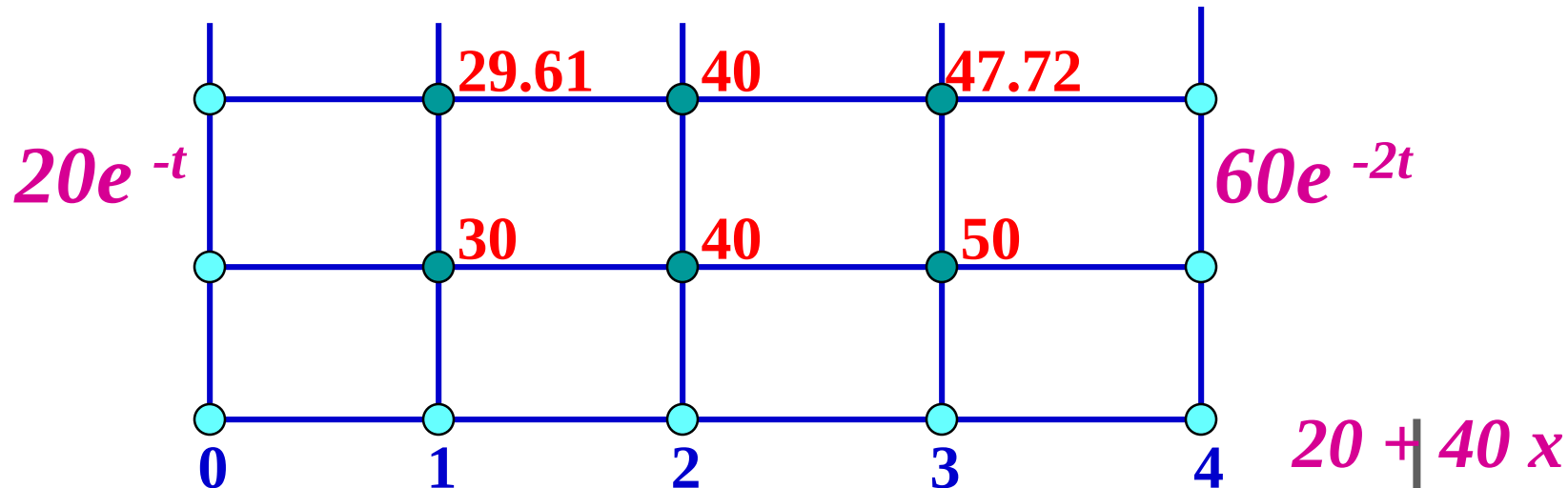
- First Step: $t = 0.05$

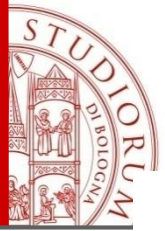
$$\left\{ \begin{aligned} u_{0,1} &= 20e^{-0.05} = 19.02458849 \\ u_{1,1} &= 0.4u_{0,0} + 0.2u_{1,0} + 0.4u_{2,0} = 0.4(20) + 0.2(30) + 0.4(40) = 30 \\ u_{2,1} &= 0.4u_{1,0} + 0.2u_{2,0} + 0.4u_{3,0} = 0.4(30) + 0.2(40) + 0.4(50) = 40 \\ u_{3,1} &= 0.4u_{2,0} + 0.2u_{3,0} + 0.4u_{4,0} = 0.4(40) + 0.2(50) + 0.4(60) = 50 \\ u_{4,1} &= 60e^{-0.10} = 54.29024508 \end{aligned} \right.$$



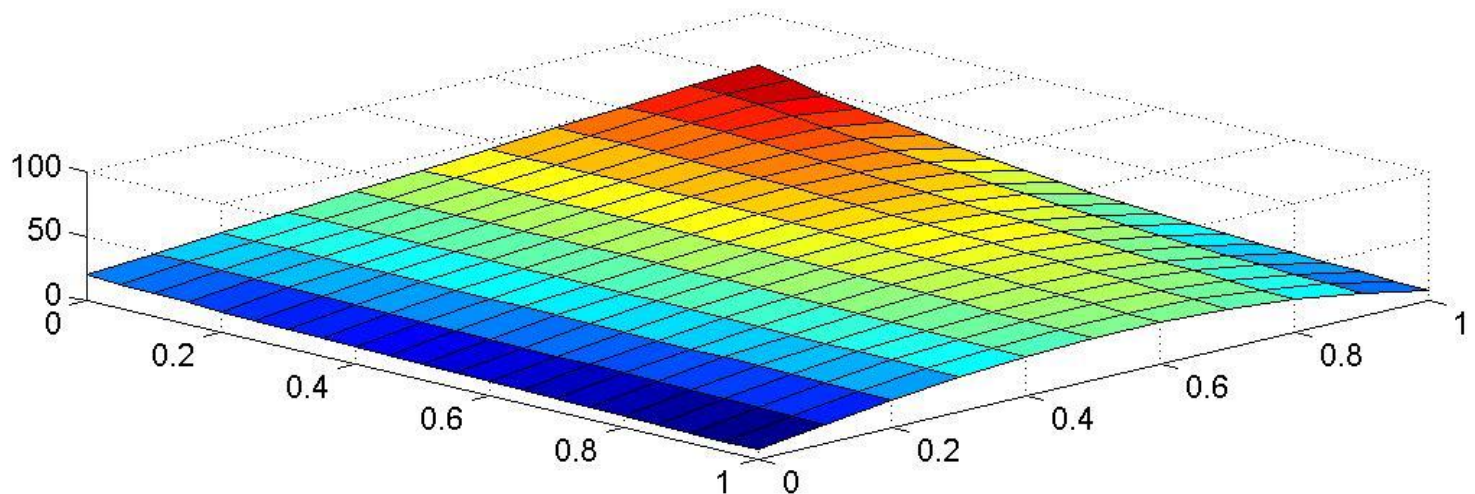
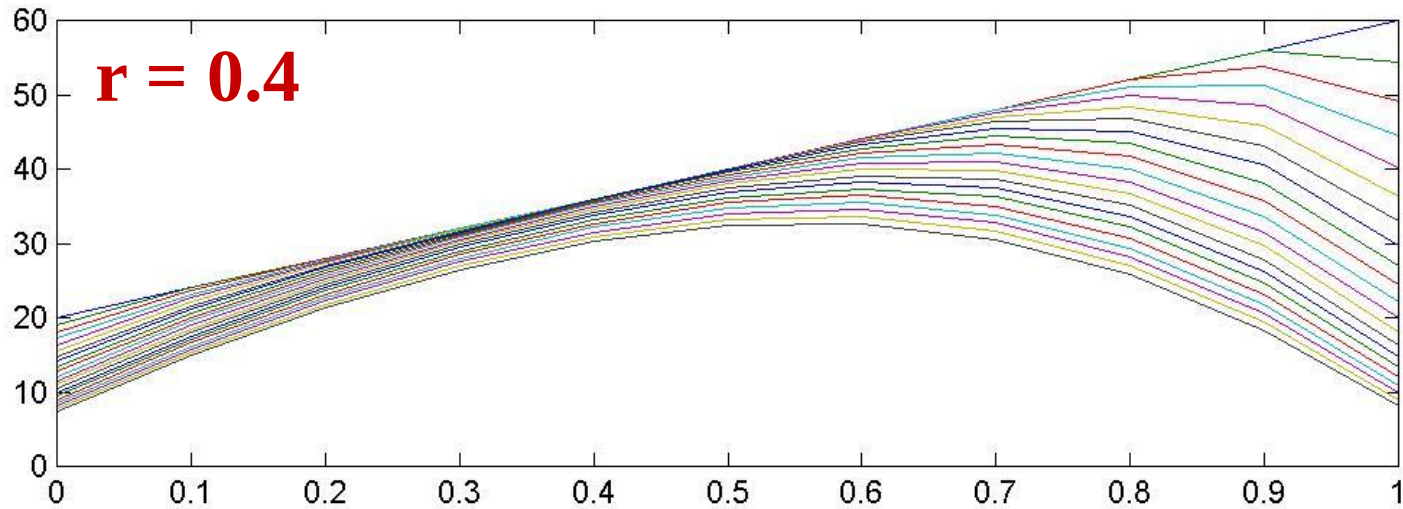
- Second step $t = 0.10$

$$\left\{ \begin{array}{l} u_{0,2} = 20e^{-0.10} = 18.09674836 \\ u_{1,2} = 0.4u_{0,1} + 0.2u_{1,1} + 0.4u_{2,1} \\ \quad = 0.4(19.02458849) + 0.2(30) + 0.4(40) = 29.6098354 \\ u_{2,2} = 0.4u_{1,1} + 0.2u_{2,1} + 0.4u_{3,1} = 0.4(30) + 0.2(40) + 0.4(50) = 40 \\ u_{3,2} = 0.4u_{2,1} + 0.2u_{3,1} + 0.4u_{4,1} \\ \quad = 0.4(40) + 0.2(50) + 0.4(54.2924508) = 47.71609803 \\ u_{4,2} = 60e^{-0.20} = 49.12384518 \end{array} \right.$$



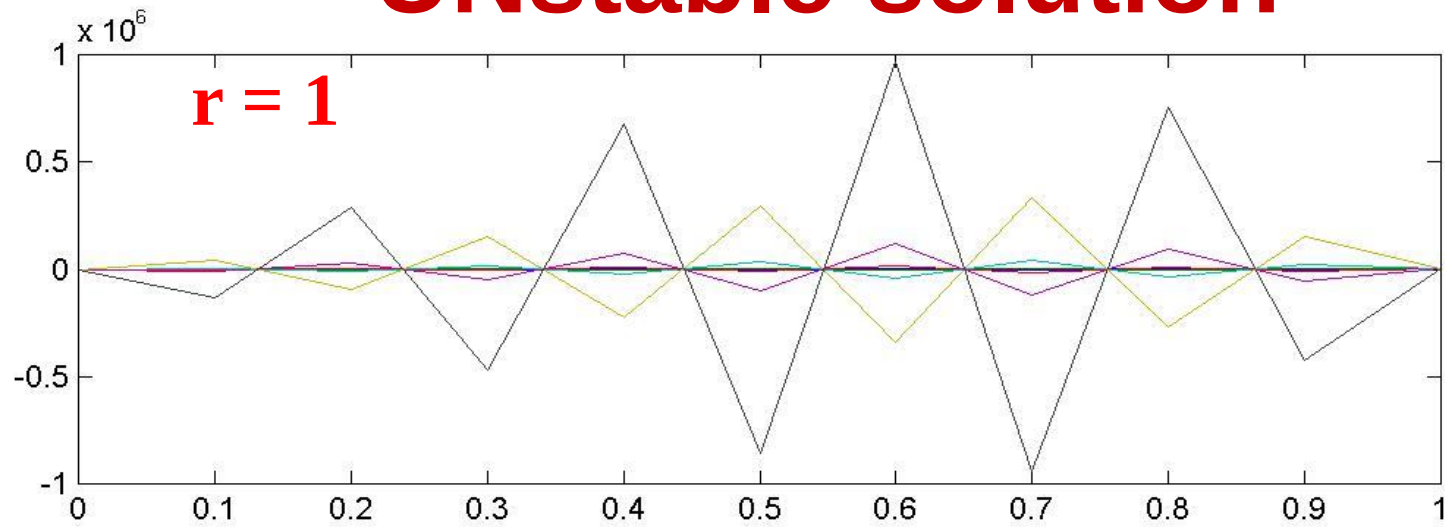


Heat Equation: BC that depend on time

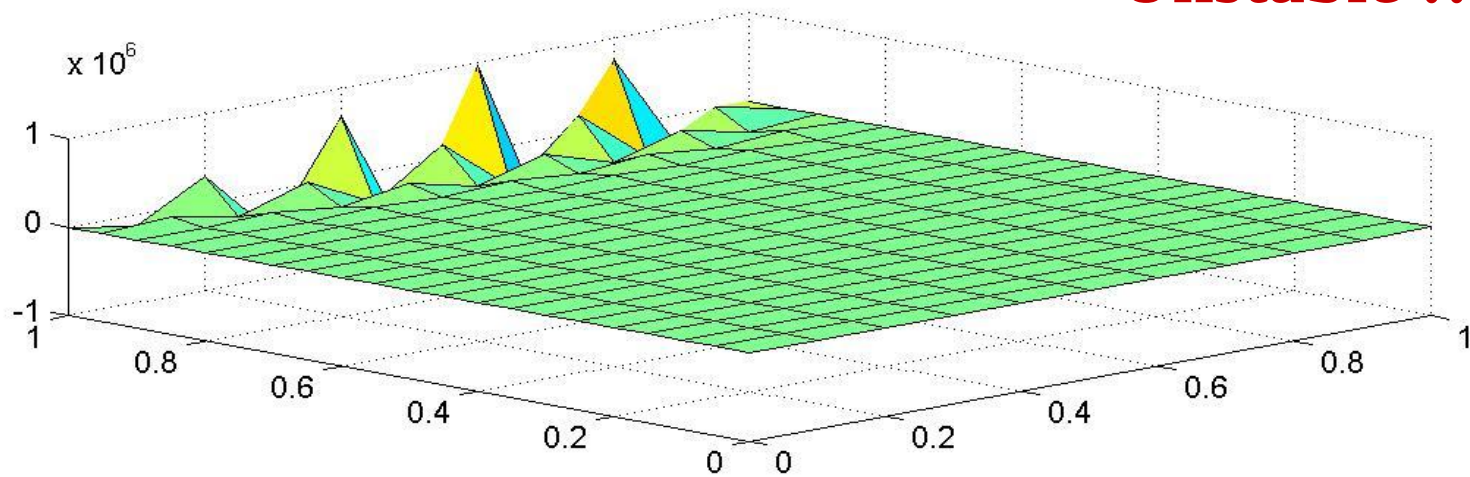


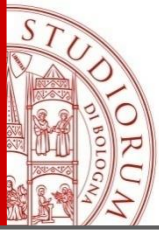


Example: explicit method, UNstable solution



Unstable !!



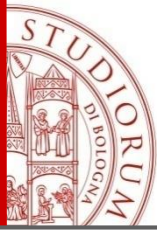


Stability Conditions

Conditioned Stability for the explicit method:

$$r(=\frac{ck}{h^2}) \leq \frac{1}{2} \quad \text{that is} \quad \Delta t \leq \frac{1}{2} \frac{\Delta x^2}{c}$$

⇒ an implicit method to avoid the instability



Implicit Method

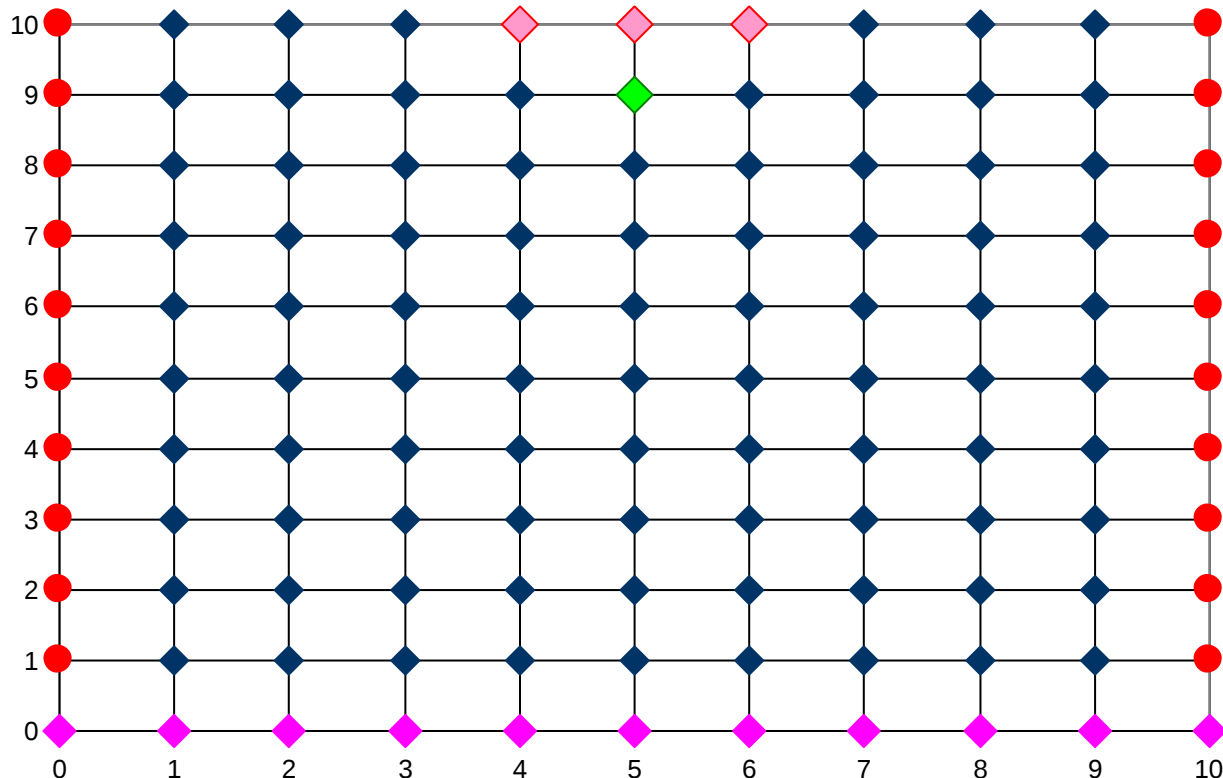
$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}; \quad 0 < x < 1, \quad 0 \leq t \leq T$$

$$\text{CI: } u(x, 0) = f(x), \quad 0 < x < 1$$

$$\text{CB: } \begin{cases} u(0, t) = g_0(t) \\ u(1, t) = g_1(t) \end{cases}, \quad 0 \leq t \leq T$$

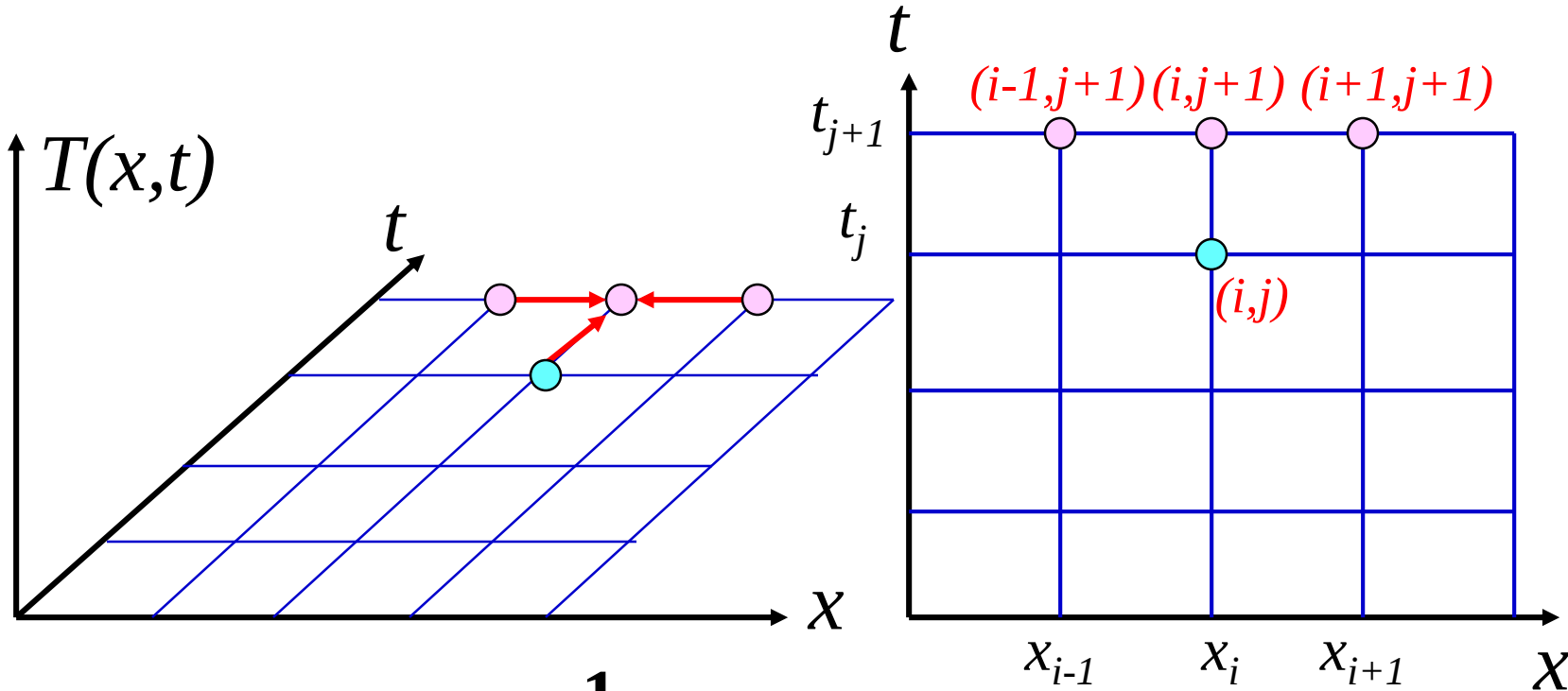
$$u(0, t) = g_0(t)$$

$$u(1, t) = g_1(t)$$



Initial Conditions : $u(x, 0) = f(x)$

Implicit Method

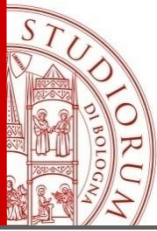


Forward difference

$$u_t = \frac{1}{k} (u_i^{j+1} - u_i^j) + O(k)$$

**Centered difference
at time $j+1$**

$$u_{xx} = \frac{1}{h^2} (u_{i-1}^{j+1} - 2u_i^{j+1} + u_{i+1}^{j+1}) + O(h^2)$$



Implicit Method

$$\frac{1}{k}(u_i^{j+1} - u_i^j) = \frac{1}{h^2}(u_{i-1}^{j+1} - 2u_i^{j+1} + u_{i+1}^{j+1})$$

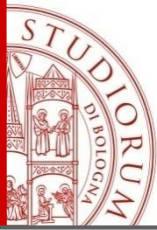
$$u_i^j = -ru_{i-1}^{j+1} + (1 + 2r)u_i^{j+1} - ru_{i+1}^{j+1}$$

Linear system to be solved at each time step t
Tridiagonal coefficient matrix (Thomas alg.)

$$\begin{bmatrix} 1+2r & -r & & & \\ -r & 1+2r & -r & & \\ & -r & 1+2r & -r & \\ & & \ddots & \ddots & -r \\ & & & -r & 1+2r \end{bmatrix} \begin{Bmatrix} u_{1,j+1} \\ u_{2,j+1} \\ u_{3,j+1} \\ \vdots \\ u_{N-1,j+1} \end{Bmatrix} = \begin{Bmatrix} u_{1,j} + ru_{0,j+1} \\ u_{2,j} \\ u_{3,j} \\ \vdots \\ u_{N-1,j} + ru_{N,j+1} \end{Bmatrix}$$

$$AU_{j+1} = U_j$$

Unconditionally stable



Stability of the Implicit Method

$$(I - r L)U_{j+1} = U_j$$

$$U_{j+1} = (I - r L)^{-1}U_j$$

$$U_{j+1} = A^{-1}U_j$$

$$U_{j+1} = A^{-2}U_{j-1} = \dots = A^{-(j+1)}U_0$$

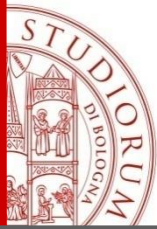
Let us introduce some perturbations E_0 on the data U_0

$$\bar{U}_0 = U_0 + E$$

$$\bar{U}_{j+1} = A^{-(j+1)}\bar{U}_0 = A^{-(j+1)}(U_0 + E) = A^{-(j+1)}U_0 + A^{-(j+1)}E$$

The effect of the error E after $j+1$ steps is

$$\|E^{(j+1)}\| \leq \|A^{-(j+1)}\|_2 \|E\|$$



Stability of the Implicit Method

$$\|E^{(j+1)}\| \leq \|A^{-(j+1)}\|_2 \|E\|$$

$$\|A^{-1}\|_2 = \rho(A^{-1}) = \max_{1 \leq p \leq N} |\lambda_p^{-1}| = \left(\min_{1 \leq p \leq N} |\lambda_p| \right)^{-1}$$

eigenvalues of L : $\lambda_p = 2(\cos(p\pi h) - 1)$ $p = 1, \dots, N-1, h = 1/N \Rightarrow$

$$\text{smallest } \lambda_1 = 2(\cos(\pi h) - 1) = 2\left(-\frac{1}{2}\pi^2 h^2 + \frac{1}{24}\pi^4 h^4 + O(h^6)\right) \approx -\pi^2 h^2 + O(h^4)$$

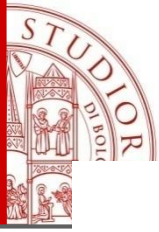
since $A = I - rL$ and $r = \frac{ck}{h^2}$

the method is

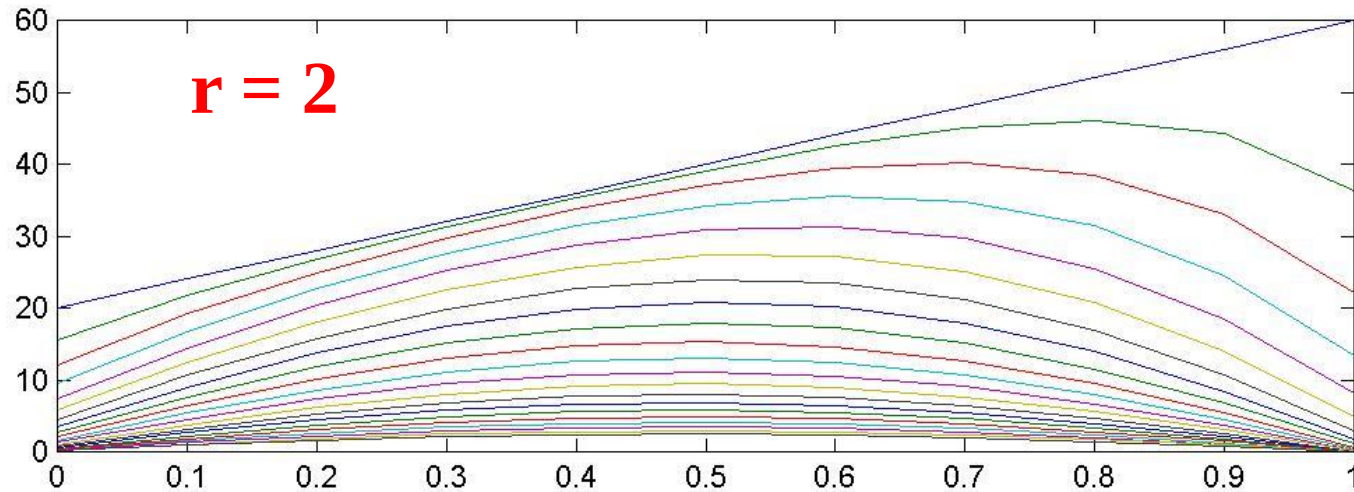
unconditionally stable

$$\|A^{-(j+1)}\|_2 \approx \frac{1}{|1 + ck\pi^2|} \leq 1$$

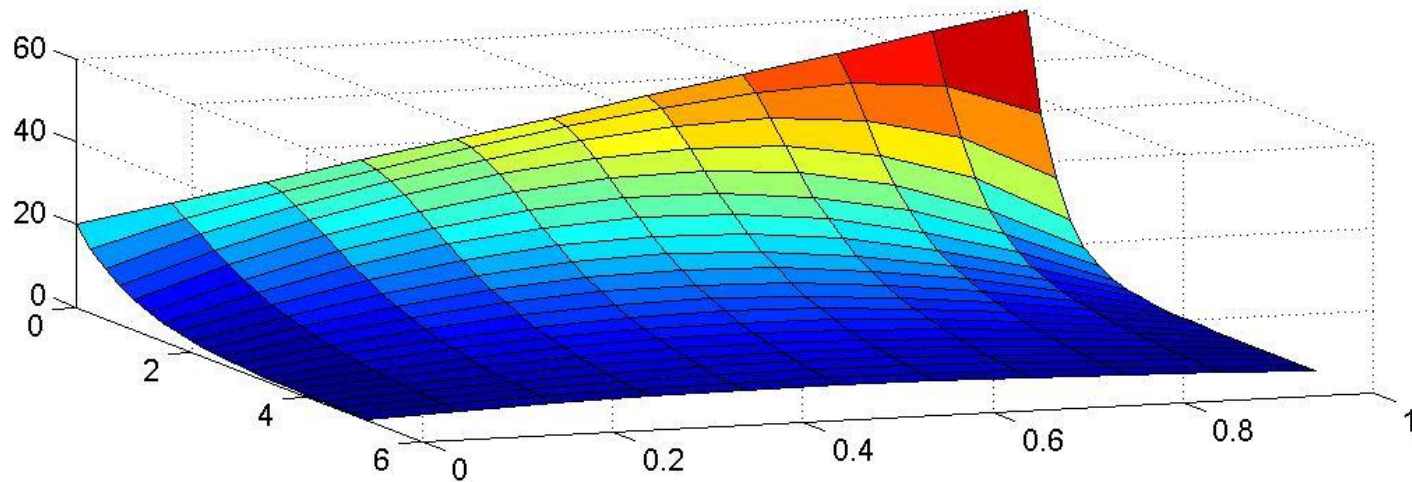
Note that $\|\cdot\| < 1$ is more restrictive than the Lax-Richtmyer constant C_T

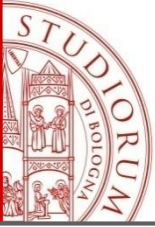


Implicit Method



Unconditionally stable first order accuracy in time





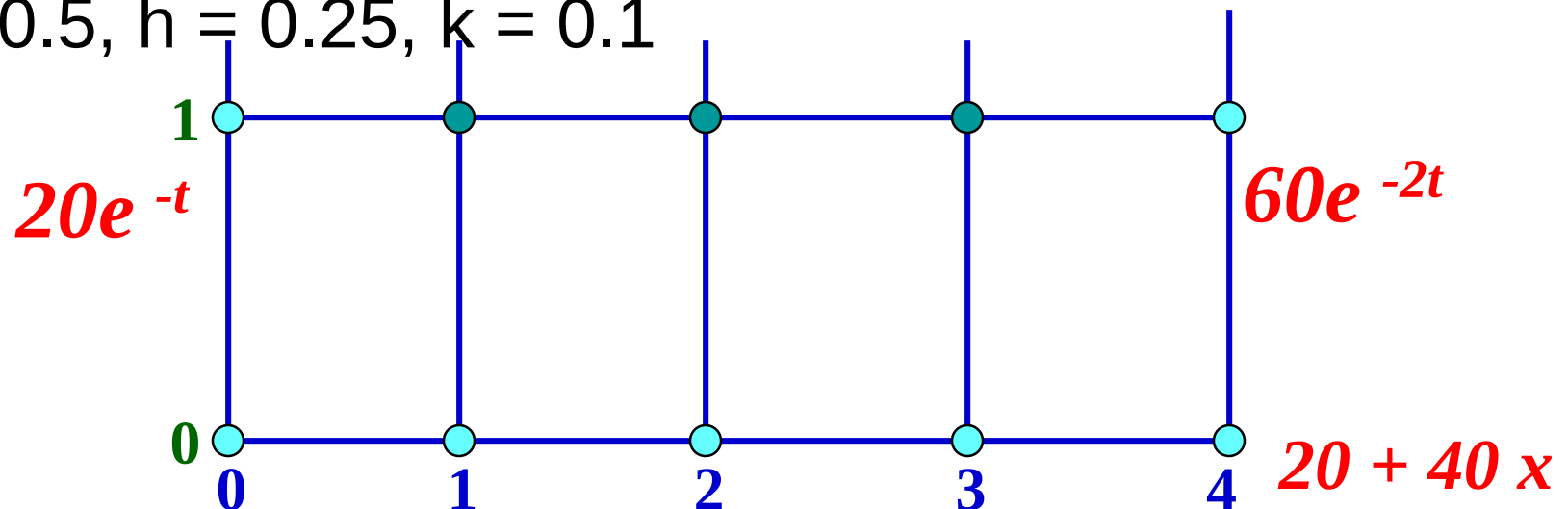
Example: Implicit Method, stable solution

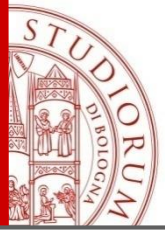
- Heat Equation (Parabolic PDE)

$$u_t = cu_{xx} ; 0 \leq x \leq 1$$

$$\begin{cases} u(x,0) = 20 + 40x \\ u(0,t) = 20e^{-t}, u(1,t) = 60e^{-2t} \end{cases}$$

$$c = 0.5, h = 0.25, k = 0.1$$





Example: Implicit Method, stable solution

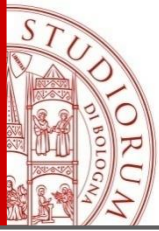
- Implicit Euler Method

$$r = \frac{ck}{h^2} = \frac{(0.5)(0.10)}{(0.25)^2} = 0.8$$

$$(-r)u_{i-1,j+1} + (1 + 2r)u_{i,j+1} + (-r)u_{i+1,j+1} = u_{i,j}$$

$$(-0.8)u_{i-1,j+1} + (2.6)u_{i,j+1} + (-0.8)u_{i+1,j+1} = u_{i,j}$$

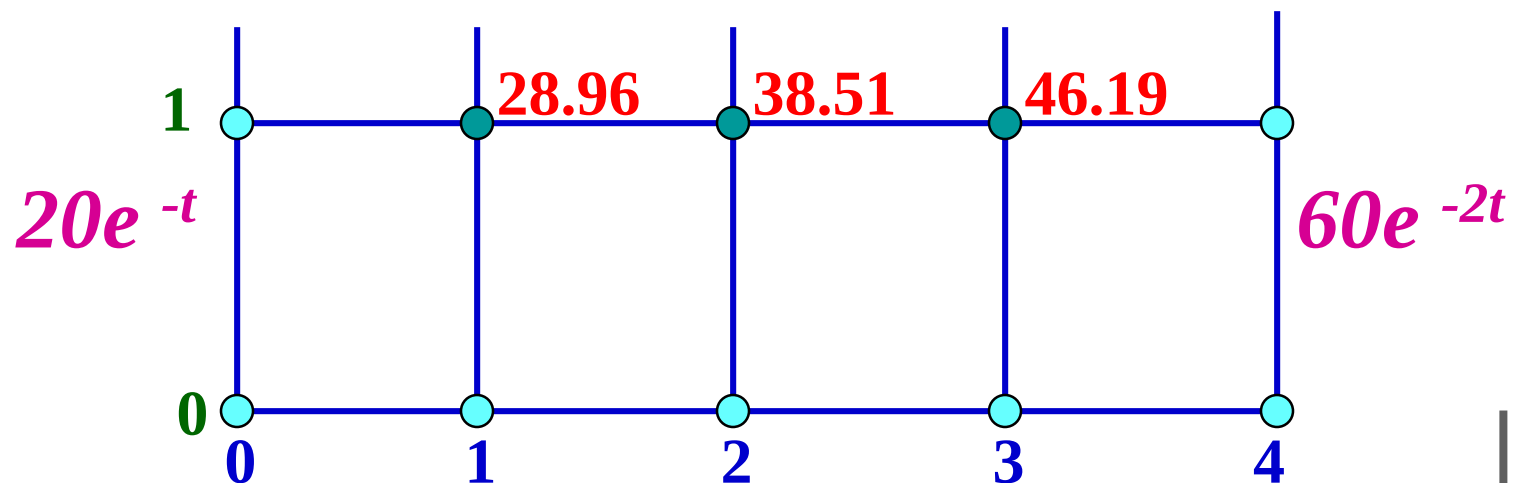
$$\begin{bmatrix} 1 + 2r & -r & 0 \\ -r & 1 + 2r & -r \\ 0 & -r & 1 + 2r \end{bmatrix} \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} u_{1,0} + ru_{0,1} \\ u_{2,0} \\ u_{3,0} + ru_{4,1} \end{Bmatrix}$$



Solve the tridiagonal system

$$\begin{bmatrix} 2.6 & -0.8 & 0 \\ -0.8 & 2.6 & -0.8 \\ 0 & -0.8 & 2.6 \end{bmatrix} \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} 30 + 0.8(20e^{-0.1}) \\ 40 \\ 50 + 0.8(60e^{-0.2}) \end{Bmatrix}$$

$$\Rightarrow \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} 28.95515793 \\ 38.50751457 \\ 46.19426454 \end{Bmatrix}$$





John Crank
1916-2006

$$u(0, t) = g_0(t)$$

Crank-Nicolson method (implicit)

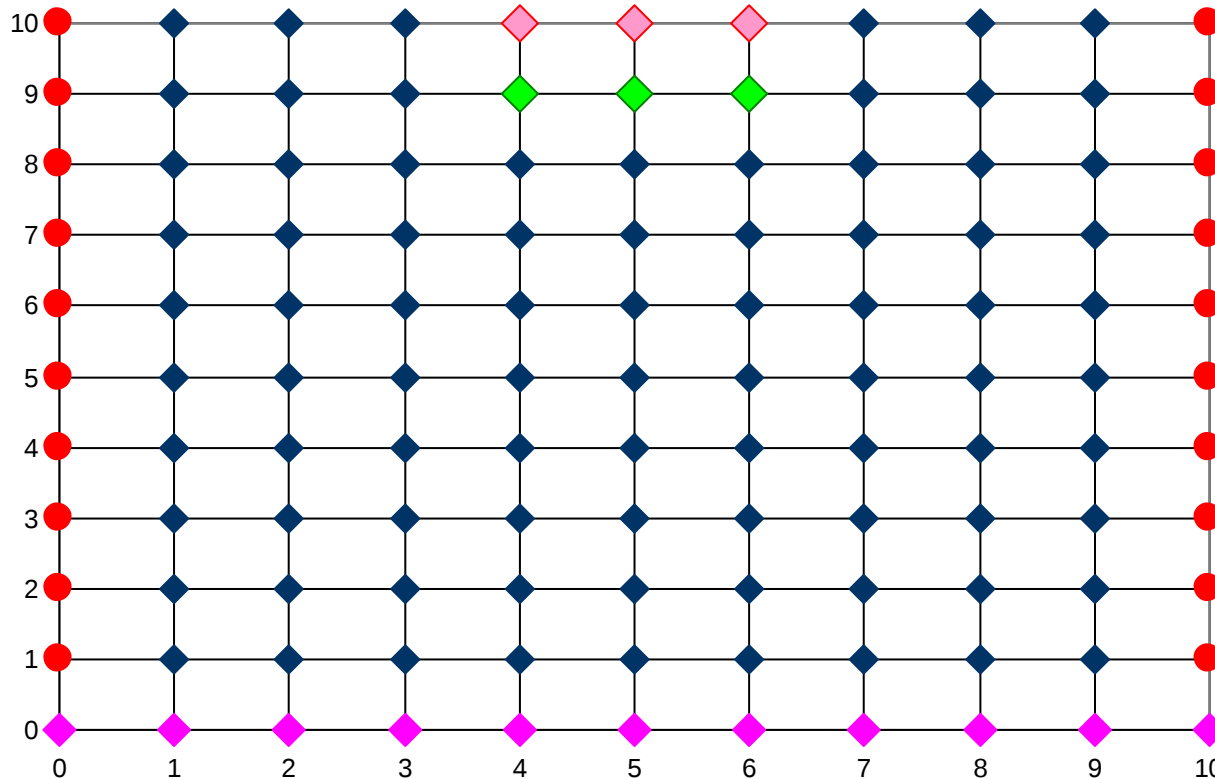
Implicit method: 1° order in time

Crank-Nicolson : 2° order in time

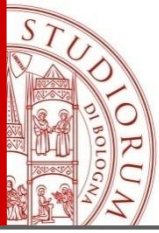


Phyllis
Nicolson
1917-1968

$$u(1, t) = g_1(t)$$



Initial Conditions: $u(x, 0) = f(x)$



Crank-Nicolson method (implicit)

- Crank-Nicolson method for heat equation

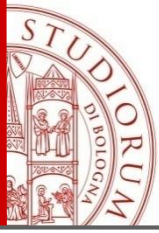
$$\frac{\partial u}{\partial t} = a \frac{\partial^2 u}{\partial x^2}; \quad 0 < x < 1, \quad 0 \leq t \leq T$$

$$u(x, 0) = f(x), \quad -\infty < x < \infty$$

- We integrate with respect to time and we use the method of trapezoids to solve the integral in the second member

$$u(x, t + \Delta t) - u(x, t) = a \int_t^{t+\Delta t} u_{xx}(x, s) ds$$

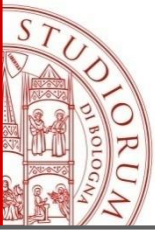
$$u(x, t + \Delta t) - u(x, t) = a \frac{\Delta t}{2} (u_{xx}(x, t) + u_{xx}(x, t + \Delta t)) \\ - \frac{1}{12} \Delta t^3 u_{xxtt}(x, \theta_t)$$



Crank-Nicolson method (implicit)

We replace the derivatives with finite differences
(results in an average of the central differences at
two successive time steps j and $j + 1$)

$$\frac{1}{k}(u_{i,j+1} - u_{i,j}) = \frac{a}{2h^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) \\ + \frac{a}{2h^2}(u_{i-1,j+1} - 2u_{i,j+1} + u_{i+1,j+1}) + O(h^2) + O(k^2)$$



Crank-Nicolson method (implicit)

$$\frac{1}{k}(u_{i,j+1} - u_{i,j}) = \frac{a}{2h^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) \\ + \frac{a}{2h^2}(u_{i-1,j+1} - 2u_{i,j+1} + u_{i+1,j+1})$$

We get the following implicit method $r = \frac{ak}{2h^2}$

$$-ru_{i-1,j+1} + (1+2r)u_{i,j+1} - ru_{i+1,j+1} = ru_{i-1,j} + (1-2r)u_{i,j} + ru_{i+1,j}$$

Linear System of N-1 linear eqs with tridiagonal coefficient matrix

$$\mathbf{A}\mathbf{U}_{j+1} = \mathbf{B}\mathbf{U}_j$$

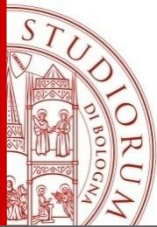
$$IC : u_{i,0} = u_0(x_i) \quad i = 0, \dots, N$$

$$BC : u(0,t) = g_0(t); \quad u(1,t) = g_1(t)$$

$$\begin{bmatrix} 1+2r & -r & 0 & & \\ -r & 1+2r & -r & & \\ & & & & \\ & & & & -r \\ & & & -r & 1+2r \end{bmatrix} \begin{bmatrix} u_{1,j+1} \\ u_{2,j+1} \\ \vdots \\ u_{i,j+1} \\ \vdots \\ u_{N-2,j+1} \\ u_{N-1,j+1} \end{bmatrix} = \begin{bmatrix} 1-2r & r & 0 & & \\ & r & 1-2r & r & \\ & & & & r \\ & & & & \\ & & & r & 1-2r \end{bmatrix} \begin{bmatrix} u_{1,j} \\ u_{2,j} \\ \vdots \\ u_{i,j} \\ \vdots \\ u_{N-2,j} \\ u_{N-1,j} \end{bmatrix} + \begin{bmatrix} +r(g_{0,j} + g_{0,j+1}) \\ 0 \\ \vdots \\ 0 \\ +r(g_{1,j} + g_{1,j+1}) \end{bmatrix}$$

$$AU_{j+1} = BU_j + v \quad A, B \in R^{(N-1) \times (N-1)}, U \in R^{N-1}$$

Truncation Error $O(k^2 + h^2)$



$$U_{j+1} = A^{-1}BU_j$$

- **Stability**

$$\|A^{-1}B\| \leq C_T \rightarrow \rho(A^{-1}B) \leq 1 \quad \text{eigenvalues of } A^{-1}B:$$

$$A^{-1}B = (I - \frac{r}{2}L)^{-1}(I + \frac{r}{2}L) \quad r = \frac{k}{h^2}$$

$$\frac{1 + \frac{k}{h^2}(\cos(p\pi h) - 1)}{1 - \frac{k}{h^2}(\cos(p\pi h) - 1)} \leq 1 \quad p = 1, \dots, N-1$$

Unconditionally stable

$$O(k^2 + h^2)$$

- **Consistency**

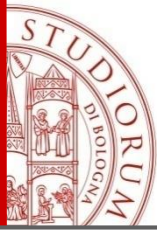
2_{nd} order accurate in h (space) and k (time)

Θ -Method

$$\frac{1}{k}(u_{i,j+1} - u_{i,j}) = \frac{c\theta}{h^2}(u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) + \frac{c(1-\theta)}{h^2}(u_{i-1,j+1} - 2u_{i,j+1} + u_{i+1,j+1})$$

Weighted Average of partial differential equations between time steps j and $j+1$.

$$\left\{ \begin{array}{ll} \theta = 0: & \text{implicit method} \\ \theta = 1: & \text{explicit method} \\ \theta = 1/2: & \text{Crank-Nicolson} \end{array} \right.$$



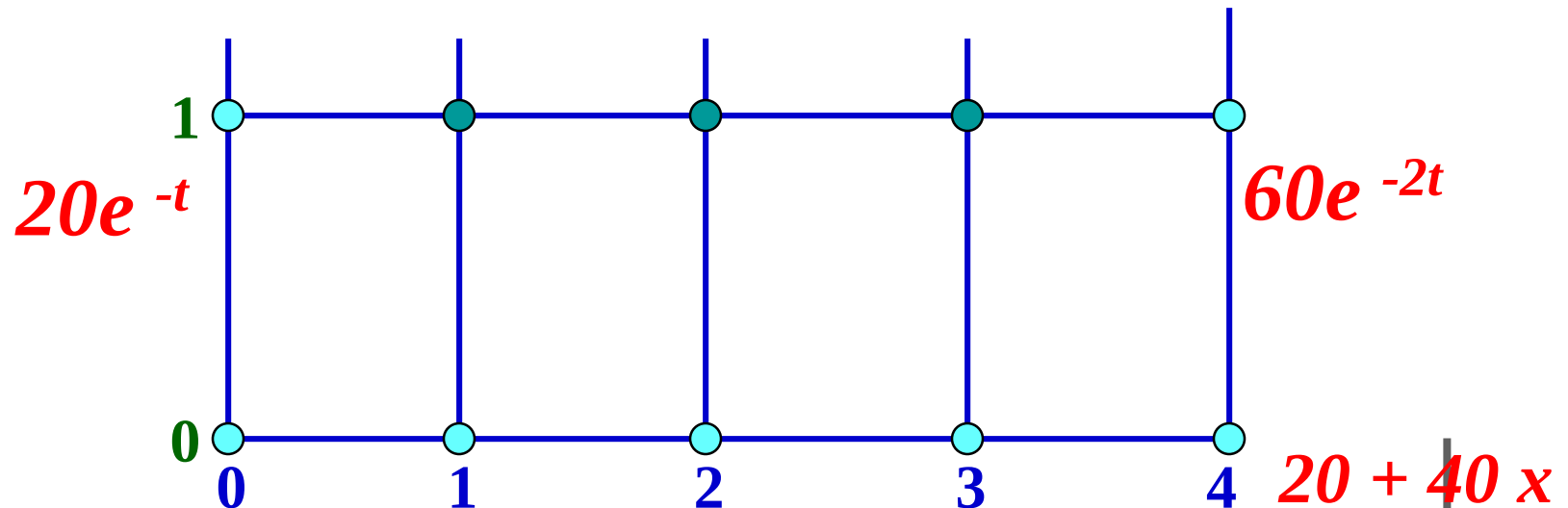
Example: CN method, stable solution

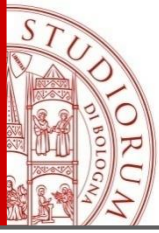
- Heat Equation (Parabolic PDE)

$$u_t = cu_{xx} ; 0 \leq x \leq 1$$

$$\begin{cases} u(x,0) = 20 + 40x \\ u(0,t) = 20e^{-t}, u(1,t) = 60e^{-2t} \end{cases}$$

$$c = 0.5, h = 0.25, k = 0.1$$





Example: CN method, stable solution

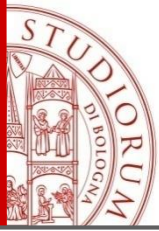
- Crank-Nicolson

$$r = \frac{ck}{h^2} = \frac{(0.5)(0.10)}{(0.25)^2} = 0.8$$

$$-\frac{r}{2}u_{i-1,j+1} + (1+r)u_{i,j+1} - \frac{r}{2}u_{i+1,j+1} = \frac{r}{2}u_{i-1,j} + (1-r)u_{i,j} + \frac{r}{2}u_{i+1,j}$$
$$-0.4u_{i-1,j+1} + 1.8u_{i,j+1} - 0.4u_{i+1,j+1} = 0.4u_{i-1,j} + 0.2u_{i,j} + 0.4u_{i+1,j}$$

- Tridiagonal matrix ($r = 0.8$)

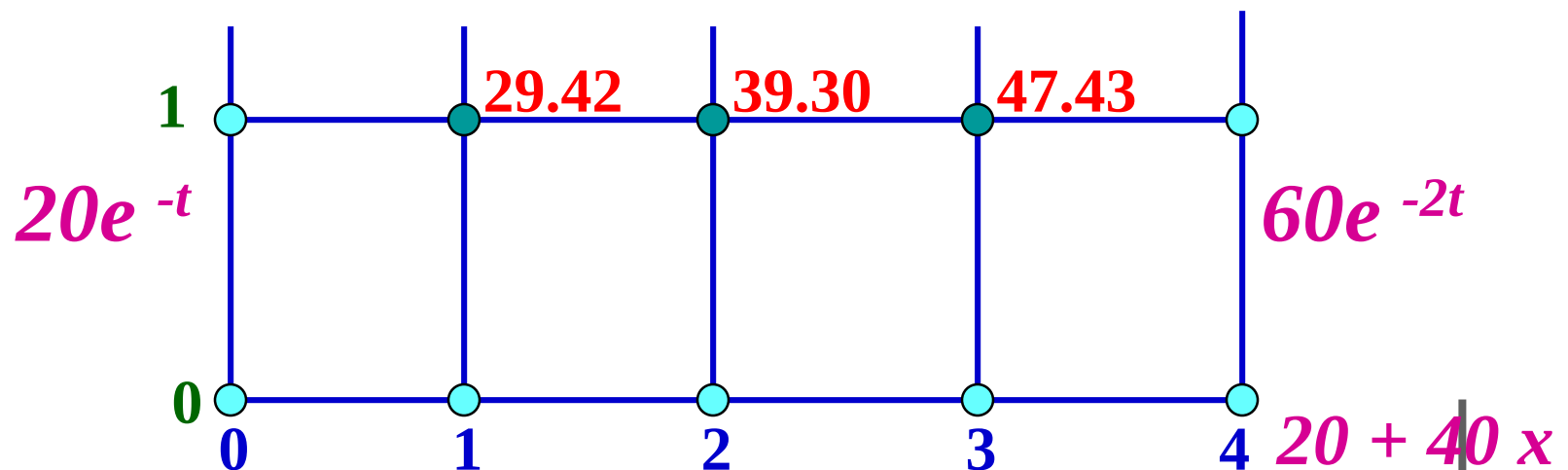
$$\begin{bmatrix} 1+r & -\frac{r}{2} & 0 \\ -\frac{r}{2} & 1+r & -\frac{r}{2} \\ 0 & -\frac{r}{2} & 1+r \end{bmatrix} \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} \frac{r}{2}u_{0,0} + (1-r)u_{1,0} + \frac{r}{2}u_{2,0} + \frac{r}{2}u_{0,1} \\ \frac{r}{2}u_{1,0} + (1-r)u_{2,0} + \frac{r}{2}u_{3,0} \\ \frac{r}{2}u_{2,0} + (1-r)u_{3,0} + \frac{r}{2}u_{4,0} + \frac{r}{2}u_{4,1} \end{Bmatrix}$$

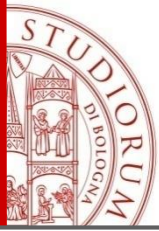


Example: CN method, stable solution

$$\begin{bmatrix} 1.8 & -0.4 & 0 \\ -0.4 & 1.8 & -0.4 \\ 0 & -0.4 & 1.8 \end{bmatrix} \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} 0.4(20) + 0.2(30) + 0.4(40) + 0.4(20e^{-0.1}) \\ 0.4(30) + 0.2(40) + 0.4(50) \\ 0.4(40) + 0.2(50) + 0.4(60) + 0.4(60e^{-0.2}) \end{Bmatrix}$$

$$= \begin{Bmatrix} 37.23869934 \\ 40 \\ 69.64953807 \end{Bmatrix} \Rightarrow \begin{Bmatrix} u_{1,1} \\ u_{2,1} \\ u_{3,1} \end{Bmatrix} = \begin{Bmatrix} 29.42144598 \\ 39.29975855 \\ 47.42746748 \end{Bmatrix}$$





Non-constant Diffusion Coefficient

$$u_t = cu_{xx}$$

c diffusion coefficient $c > 0$, in general, $c(x, t) > 0$ depends on x and t

Explicit

$$u_{i,j+1} = u_{i,j} + \frac{c_{i,j}k}{h^2} (u_{i-1,j} - 2u_{i,j} + u_{i+1,j})$$

$$k \leq \frac{c_{i,j}}{h^2}, \quad k \leq \frac{\bar{c}_j}{h^2} \quad \bar{c} = \max_i c_{i,j}$$

Implicit

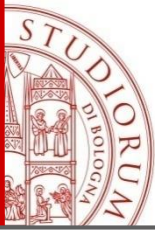
$$u_{i,j+1} = u_{i,j} + \frac{c_{i,j+1}k}{h^2} (u_{i-1,j+1} - 2u_{i,j+1} + u_{i+1,j+1})$$

Θ -Method

$$c_{i,j+1/2} = (c_{i,j} + c_{i,j+1}) / 2$$

or

$$\frac{1}{k} (u_{i,j+1} - u_{i,j}) = \frac{c_{i,j}\theta}{h^2} (u_{i-1,j} - 2u_{i,j} + u_{i+1,j}) + \frac{c_{i,j+1}(1-\theta)}{h^2} (u_{i-1,j+1} - 2u_{i,j+1} + u_{i+1,j+1})$$



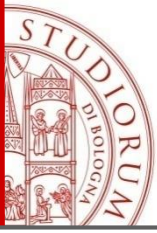
Method of lines (MOL)

- We first discretize the PDE in space alone, which gives a large system of ODEs with each component of the system corresponding to the solution at some grid point, as a function of time. (**semi-discretization**)
- Then we solve the system using one of the methods for ODEs that we have previously studied:
Explicit Euler, Implicit Euler, Θ -method,...).
- Finally we fix the boundary conditions.

$$u_t - au_{xx} = 0 \quad + IC \quad + BC \quad (i.e. periodic)$$

$$\text{Let } U(t) = [u_0(t), u_1(t), \dots, u_n(t)]$$

$$\rightarrow U'(t) = \frac{a}{h^2} LU(t) \quad \text{System of ODEs}$$

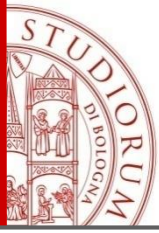


Method of lines (MOL)

$$\begin{bmatrix} u_0'(t) \\ u_1'(t) \\ \ddots \\ u_i'(t) \\ \ddots \\ u_{N-1}'(t) \\ u_N'(t) \end{bmatrix} = \frac{a}{h^2} \begin{bmatrix} -2 & 1 & 0 & & \\ 1 & -2 & 1 & & \\ & \ddots & \ddots & 1 & \\ & & 1 & -2 & \\ & & & & \ddots \end{bmatrix} \begin{bmatrix} u_0(t) \\ u_1(t) \\ \ddots \\ u_i(t) \\ \ddots \\ u_{N-1}(t) \\ u_N(t) \end{bmatrix}$$

$$U'(t) = \frac{a}{h^2} L U(t) \quad L \in \mathbb{R}^{(N+1) \times (N+1)}, U \in \mathbb{R}^{N+1}$$

L discretization of the Laplace operator



Method of lines (MOL): apply discretization in time

Explicit Euler Method $U^{n+1} = (I + r L)U^n \quad r = \frac{ak}{h^2}$

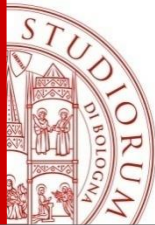
Implicit Euler Method $(I - r L)U^{n+1} = U^n$

Large Linear System of equations

Local Truncation Error $O(k + h^2)$

Crank-Nicolson Method $(I - \frac{r}{2} L)U^{n+1} = (I + \frac{r}{2} L)U^n$

Local Truncation Error $O(k^2 + h^2)$



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